

1 Dynamically and Kinematically Consistent Global Ocean
2 Circulation and Ice State Estimates

3 DRAFT

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Abstract

The World Ocean Circulation Experiment (WOCE) drove the development of estimates of the decadal scale time evolving general circulation that are dynamically and kinematically consistent. A long time-scale, and a goal of estimation rather than prediction, preclude the use of meteorological methods called “data assimilation (DA).” Instead, “state estimation” methods are reviewed here and distinguished from DA. Results from the dynamically-consistent ECCO family of solutions based upon least-squares Lagrange multipliers (adjoints) are used to discuss the determination of the dominant elements of the circulation in the period since 1992—which marked the beginning of the satellite altimetric record.

1 Introduction

The goal of what we call “state estimates” of the oceans arose directly out of the plans for the World Ocean Circulation Experiment (WOCE). That program, out of necessity, employed in a pragmatic way observational tools of a very wide diversity of type—including classical hydrography, current meters, tracers, satellite altimeters, floats, and drifters. The designers of WOCE realized that to obtain a coherent picture of the global ocean circulation approaching a time-scale of a decade, they would require some form of synthesis method: one capable of combining very disparate observational types, but also having greatly differing space-time sampling, and geographical coverage.

Numerical weather forecasting, in the form of what had become known as “data assimilation” (DA), was a known analogue of what was required: a collection of tools for combining the best available global numerical model representation of the ocean with any and all data, suitably weighted to account for both model and data errors (e.g., Talagrand, 1997; Kalnay, 2003; Evensen, 2009). Several major, and sometimes ignored, obstacles existed in employing meteorological methods for the oceanic problem. These included the large infrastructure used to carry out DA within the national weather forecast centers—organizations for which no oceanographic equivalent existed or exists. DA had developed for the purposes of *forecasting* over time scales of hours to a few days, whereas the climate goals of WOCE were directed at time scales of years to decades with a goal of understanding and *not* forecasting. Another, more subtle, difficulty was the WOCE need for state estimates capable of being used for global-scale energy, heat, and water cycle budgets. Closed global budgets are of little concern to a weather forecaster—as their violation has no impact on short-range prediction skill, but they are crucial to the understanding of climate change. Construction of closed budgets is also rendered physically impossible by the forecasting goal: solutions “jump” towards the data at every analysis time, usually every six hours, introducing spurious sources and sinks of basic properties.

55 Because of these concerns, the widespread misunderstanding of what DA usually does, and
56 what oceanographers actually require, the first part of this essay is devoted to a sketch of
57 the basic principles of DA, and the contrast with methods required in practice for use for
58 climate-relevant state estimates. More elaborate accounts can be found in Wunsch (2006),
59 and Wunsch and Heimbach (2007) among others. Within the meteorological literature itself,
60 numerous publications exist (e.g., Trenberth et al., 1995, 2001; Bengtsson et al., 2004; Bromwich
61 and Fogt, 2004; Bromwich et al., 2004, 2007, 2011; Nicolas and Bromwich, 2011; Thorne, 2008)
62 warning against the use of DA and the associated “reanalyses” for the study of climate change.
63 These warnings have been widely disregarded.

64 A theme of this Chapter is that both DA and state estimation can be understood from
65 elementary principles, ones not going beyond beginning calculus. Those concepts must be dis-
66 tinguished from the far more difficult numerical engineering problem of finding practical methods
67 capable of coping with large volumes of data, large model state dimensions, and a variety of
68 computer architectures. But one can understand and use an automobile without being expert
69 in the manufacture of an internal combustion engine or of the chemistry of tire production.

70 At the time of the writing of the first WOCE volume, Siedler et al. (2001), two types
71 of large-scale synthesis existed: (1) the time-mean global inverse results of Macdonald (1998)
72 based upon the pre-WOCE hydrography and that of Ganachaud (2003b) using the WOCE
73 hydrographic sections. (2) Preliminary results from the first ECCO (Estimating the Circulation
74 and Climate of the Ocean) synthesis (Stammer et al., 2002) were based upon a few years of data
75 and comparatively coarse resolution models. Talley et al. (2001) summarized these estimates,
76 but little time had been available for their digestion.

77 In the intervening years, Lumpkin and Speer (2007) produced a revision of the Ganachaud
78 results using somewhat different assumptions, but with similar results, and a handful of other
79 static global estimates (e.g., Schlitzer, 2007) appeared. The ECCO project greatly extended
80 its capabilities and duration for time-dependent estimates. A number of regional, assumed
81 steady-state, box inversions also exist (e.g., Macdonald et al., 2009).

82 As part of his box inversions, Ganachaud (2003a) had shown that the dominant errors in
83 trans-oceanic property transports of volume (mass), heat (enthalpy), salt, etc. arose from the
84 temporal variability. Direct confirmation of that inference can be seen in the ECCO-based
85 time-varying solutions, and from in situ measurements (Rayner et al., 2011). So-called synoptic
86 sections spanning ocean basins and which had been the basis for most global circulation pictures
87 at best produce “blurred” snapshots of transport properties. We are now well-past the time in
88 which they can be labelled and interpreted as being the time-average. A major result of WOCE
89 was to confirm the conviction that *the ocean must be observed and treated as a fundamentally*

90 *time-varying system*, especially for any property involving the flow field. Gross scalar proper-
 91 ties such as the temperature or nitrogen concentrations have long been known to be stable on
 92 the largest scales: that their distributions are nonetheless often dominated by intense temporal
 93 fluctuations, sometimes involving very high wavenumbers, represents a major change in under-
 94 standing of the classical ocean properties. That understanding inevitably drives one towards
 95 state estimation methods.

96 2 Definition

97 Consider any model of a physical system satisfying known equations, written generically in
 98 discrete time as,

$$\mathbf{x}(t) = \mathbf{L}(\mathbf{x}(t - \Delta t), \mathbf{q}(t - \Delta t), \mathbf{u}(t - \Delta t)), \quad 1 \leq t \leq t_f = M\Delta t, \quad (1) \quad \{\text{model1}\}$$

99 where $\mathbf{x}(t)$ is the “state” at time t , discrete at intervals Δt , and includes those prognostic
 100 or dependent variables usually computed by a model, such as temperature or salinity in an
 101 advection-diffusion equation or a stream function in a flow problem. $\mathbf{q}(t)$ denotes known forc-
 102 ings, sources, sinks, boundary and initial conditions, and internal model parameters, and $\mathbf{u}(t)$
 103 is any such elements that are regarded as only partly or wholly unknown, hence subject to ad-
 104 justment and termed independent or control variables (or simply “controls”). \mathbf{L} is an operator
 105 and can involve a large range of calculations, including derivatives, or integrals or any other
 106 mathematically defined function. In practice, it is usually a computer code working on arrays
 107 of numbers. (Notation is approximately that of Wunsch, 2006.) Time, $t = m\Delta t$, is assumed
 108 to be discrete, with $m = 0, \dots, M$, as that is almost always true of models run on computers.¹
 109 Note that the steady-state situation is a special case, in which one writes an additional rela-
 110 tionship, $\mathbf{x}(t) = \mathbf{x}(t - \Delta t)$ and \mathbf{q}, \mathbf{u} are then time-independent. For computational efficiency,
 111 steady models are normally rewritten so that time does not appear at all, but that step is not
 112 necessary. Thus the static box inverse methods and their relatives such as the beta-spiral are
 113 special cases of the ocean estimation problem.

114 Useful observations at time t are all functions of the state and, in almost all practical situa-
 115 tions, are a linear combination of one or more state vector elements,

$$\mathbf{y}(t) = \mathbf{E}(t) \mathbf{x}(t) + \mathbf{n}(t), \quad 0 \leq t \leq t_f, \quad (2) \quad \{\text{data1}\}$$

¹An interesting mathematical literature surrounds state estimation carried out in continuous time and space in formally infinite dimensional spaces. Most of it proves irrelevant for calculations on computers which are always finite dimensional. Digression into functional analysis can be needlessly distracting.

116 where $\mathbf{n}(t)$ is the inevitable noise in the observations and $t_f = M\Delta t$. $\mathbf{y}(t)$ is a vector of
 117 whatever observations of whatever, diverse, type are available at t . (Uncertain initial conditions
 118 are included here at $t = 0$, representing them as noisy observations.) Standard matrix-vector
 119 notation is being used. In a steady-state formulation, parameter t would be suppressed. (On
 120 rare occasions, data are a nonlinear combination of the state vector: an example would be a
 121 speed measurement in terms of two components of the velocity, or a frequency spectrum for
 122 some variable is known. Methods exist, not discussed here, for dealing with such observations.
 123 Observations relating to the control vector may exist, and one easy approach to using them is
 124 to redefine elements of $\mathbf{u}(t)$ as being part of the state vector.) The “state estimation problem”²
 125 is defined as determining $\tilde{\mathbf{x}}(t)$, $0 \leq t \leq t_f$, $\tilde{\mathbf{u}}(t)$, $0 \leq t \leq t_f - \Delta t$, exactly satisfying both Eqs.
 126 (1), and (2). Tildes here denote estimates to distinguish them from the true values.

127 *Important Note:* “exact” satisfaction of Eq. (1) must be understood as meaning the model
 128 *after adjustment* by $\tilde{\mathbf{u}}(t)$. Because $\mathbf{u}(t)$ can represent, if necessary, very complex, nonlinear,
 129 and large changes to the original model, which is usually defined with $\mathbf{u}(t) = \mathbf{0}$, the adjusted
 130 model can be very different from the initial version. *But the adjusted model is known, fully*
 131 *specified, and exactly satisfied*, and is what is used for discussion of the physics or chemistry,
 132 etc. It thus differs in a fundamental way from other types of estimate rendered discontinuous
 133 by “data injection,” or forcing to data, during the final forward calculation.

Typically, one must also have some knowledge of the statistics of the controls, $\mathbf{u}(t)$, and
 observation noise, $\mathbf{n}(t)$, commonly as the first and second-order moments,

$$\langle \mathbf{u}(t) \rangle = \mathbf{0}, \quad \langle \mathbf{u}(t) \mathbf{u}(t')^T \rangle = \mathbf{Q}(t) \delta_{tt'} \quad 0 \leq t \leq t_f - \Delta t = (M - 1) \Delta t, \quad (3a) \quad \{\text{stat1}\}$$

$$\langle \mathbf{n}(t) \rangle = \mathbf{0}, \quad \langle \mathbf{n}(t) \mathbf{n}(t')^T \rangle = \mathbf{R}(t) \delta_{tt'} \quad 0 \leq t \leq t_f = M\Delta t \quad (3b) \quad \{\text{stat2}\}$$

134 The brackets denote expected values and superscript T is the vector or matrix transpose.

135 In generic terms, the problem is one of *constrained estimation/optimization*, in which, usually,
 136 one seeks to minimize both the normalized quadratic model-data differences,

$$\left\langle (\mathbf{y}(t) - \mathbf{E}(t) \mathbf{x}(t))^T \mathbf{R}^{-1}(t) (\mathbf{y}(t) - \mathbf{E}(t) \mathbf{x}(t)) \right\rangle \quad (4)$$

137 and the normalized independent variables (“controls”),

$$\left\langle \mathbf{u}(t)^T \mathbf{Q}^{-1}(t) \mathbf{u}(t) \right\rangle. \quad (5)$$

138 — subject to the exact satisfaction of the *adjusted model* in Eq. (1).

²A terminology borrowed from control theory (e.g., Gelb, 1974).

For data sets and controls that are Gaussian or nearly so, the problem as stated is equivalent to weighted least-squares minimization of the scalar,

$$J = \sum_{m=0}^M (\mathbf{y}(t) - \mathbf{E}(t) \tilde{\mathbf{x}}(t))^T \mathbf{R}^{-1}(t) (\mathbf{y}(t) - \mathbf{E}(t) \tilde{\mathbf{x}}(t)) + \sum_{m=0}^{M-1} \tilde{\mathbf{u}}(t)^T \mathbf{Q}^{-1}(t) \tilde{\mathbf{u}}(t), \quad t = m\Delta t, \quad (6) \quad \{\text{data3}\}$$

139 subject to Eq. (1). It is a *PDE-constrained least-squares problem*, and nonlinear if the model
 140 or the observations are nonlinear. The uncertain initial conditions, contained implicitly in Eq.
 141 (6), are readily written out separately if desired.

142 In comparing the solutions to DA, note that *the latter problem is different*. It seeks to
 143 minimize,

$$\text{diag} \left\langle (\tilde{\mathbf{x}}(t_{\text{now}} + \tau) - \mathbf{x}(t_{\text{now}} + \tau)) (\tilde{\mathbf{x}}(t_{\text{now}} + \tau) - \mathbf{x}(t_{\text{now}} + \tau))^T \right\rangle, \quad (7) \quad \{\text{varmin}\}$$

144 that is the variance of the state about the true value at some time *future* to t_{now} . Brackets again
 145 denote the expected value. The role of the model is to make the forecast, by setting $\mathbf{u}(t) = 0$,
 146 $t_{\text{now}} + \Delta t \leq t \leq t_{\text{now}} + \tau$, because it is unknown, and starting with the most recent estimate
 147 $\tilde{\mathbf{x}}(t_{\text{now}})$ at t_{now} . Eq. (7) is itself equivalent to a requirement of minimum square deviation at
 148 $t_{\text{now}} + \tau$. A bit more will be said about this relationship below.

149 Model error deserves an extended discussion by itself. A consequence of exact satisfaction of
 150 the model equations is that we assume the discretized version of Eq. (1) to be error-free. Model
 151 errors comes in roughly three flavors: (a) the equations are incomplete or an approximated form
 152 of the real system; (b) errors are incurred in their discretization (e.g., numerical diffusion); (c)
 153 sub-grid scale parameterizations are incomplete, and/or their parameter choices sub-optimal.
 154 Methods exist to quantify these errors in an estimation framework. For example, an explicit
 155 error term may be introduced in Eq. (1) and whose estimation would become part of the least-
 156 squares optimization. Problems arise in practice when observational coverage is insufficient to
 157 achieve adequate partition of errors between those in the explicit error terms and those in the
 158 initial and boundary conditions. Furthermore, the error terms are effectively source terms in
 159 the tendency equations that violate dynamic and kinematic consistency. The approach taken in
 160 ECCO is to move toward estimating three-dimensional fields for the most important parameters
 161 as part of the gradient-based optimization (e.g., Ferreira et al., 2005; Stammer, 2005), thus
 162 rendering the problem one of combined state and parameter estimation. Adjustments required
 163 to compensate for model errors may be projected into the parameter estimates.

164 Most of the fundamental principles of practical state estimation and of DA can be understood
 165 from the common school problem of the least-squares fitting of lines and curves to data in one

166 dimension. The central point is that the *concepts* of state estimation and data assimilation are
 167 very simple; but it is equally simple to surround them with an aura of mystery and complexity
 168 that is needless for anyone who wishes primarily to understand the meaning of the results.

169 3 Data Assimilation and the Reanalyses

Despite the technical complexities of the numerical engineering practice, DA and what are called “reanalyses” should be understood as approximate methods for obtaining a solution of a least-squares problem. Using the same notation as in Eq. (7), consider again an analysis time, $t_{now} = t_{ana} + \tau$, when data have become available, and where t_{ana} is the previous analysis time, $\tau > 0$, typically 6 hours earlier. The weather forecaster’s model has been run forward to make a prediction, $\tilde{\mathbf{x}}(t_{now}, -)$, with the minus sign denoting that newer observations have *not yet* been used. The new observations are $\mathbf{E}(t_{now}) \mathbf{x}(t_{now}) + \mathbf{n}(t_{now}) = \mathbf{y}(t_{now})$. With some understanding of the quality of the forecast, expressed in the form of an uncertainty matrix (2nd moments about the truth) called $\mathbf{P}(t_{now}, -)$, and of the covariance matrix of the observational noise, $\mathbf{R}(t_{now})$, the best combination in the L_2 -norm of the information of the model and the data is the minimum of,

$$J_1 = (\tilde{\mathbf{x}}(t_{now}) - \tilde{\mathbf{x}}(t_{now}, -))^T \mathbf{P}(t_{now}, -)^{-1} (\tilde{\mathbf{x}}(t_{now}) - \tilde{\mathbf{x}}(t_{now}, -)) + \quad (8) \quad \{\text{kalm2}\}$$

$$(\mathbf{y}(t_{now}) - \mathbf{E}(t_{now}) \mathbf{x}(t_{now}))^T \mathbf{R}(t_{now})^{-1} (\mathbf{y}(t_{now}) - \mathbf{E}(t_{now}) \mathbf{x}(t_{now})),$$

170 and whose least-squares minimum for a linear model is given rigorously by the Kalman filter. In
 171 DA practice, only some very rough approximation to that minimum is sought and obtained. True
 172 Kalman filters are never used for prediction in real geophysical fluid flow problems as they are
 173 computationally overwhelming (for more detail, see e.g., Wunsch, 2006). *Notice that J_1 assumes*
 174 *that a summation of errors is appropriate, even in the presence of strong nonlinearities.*

175 A brief excursion into meteorological “reanalyses” is worthwhile here for several reasons:
 176 (1) They are often used as an atmospheric “truth” to drive ocean, ice, chemical, and biological
 177 models. (2) A number of ocean circulation estimates have followed their numerical engineering
 178 methodology. (3) With the long history of the atmospheric data assimilation effort, many have
 179 been unwilling to believe that any alternatives exist.

180 Note that the “analysis” consists of an operational weather model run in conventional pre-
 181 diction mode, analogous to the simple form described in the previous section, adjusted, and
 182 thus displaying discontinuities at the analysis times, by attempts to approximately minimize
 183 J_1 . Because of the operational/real-time requirements, only a fraction of the global operational
 184 meteorological observations are relayed and quality-controlled in time to be available at the

185 time of analysis. Furthermore, because models have changed so much over the years, the stored
186 analyses are inhomogeneous in the underlying physics³ and model codes. Oceanographers have
187 no such products at this time; “*analyses*” in the meteorological sense do not exist, and thus the
188 jargon “reanalysis” for ocean state estimates is inappropriate.

189 Meteorological reanalysis is the recomputation, using the same prediction methodology as
190 previously used in the analysis, but with the differences that (1) the model code and combination
191 methodology are held fixed over the complete time duration of the calculation (e.g., over 50 years)
192 thus eliminating artificial changes in the state from model or method improvements and, (2)
193 including many data that arrived too late to be incorporated into the real-time analysis (see
194 Kalnay, 2003; Evensen, 2009).

195 Estimated states still have the same discontinuities at the analysis times when the model is
196 forced towards the data. Of even greater significance for oceanographic and climatic studies
197 are the temporal shifts induced in the estimates by the major changes that have taken place in
198 the observational system over several decades—most notably, but not solely, the appearance of
199 meteorological satellites. Finally, no use is made of the information content in the observations
200 of the *future* evolution of the state.

201 Although as already noted above, clear warnings have appeared in the literature—that spu-
202 rious trends and values are artifacts of changing observation systems (and see for example,
203 Marshall, et al., 2002; Elliott and Gaffen, 1992; Thompson et al., 2008)—the reanalyses are
204 rarely used appropriately, meaning with the recognition that they are subject to large errors.
205 In Fig. 1 for example, the jump in precipitation minus evaporation ($P - E$) with the advent
206 of the polar orbiting satellites implies either that the unspecified error estimates prior to that
207 time must, at a minimum, encompass the jump, and/or that computation has been erroneous,
208 or that a remarkable coincidence has occurred. But even the smaller transitions in $P - E$, e.g.,
209 over the more recent period 1992 onward, are likely too large to be physical; see Table 1.

210 Figure 2 and other, similar ones, are further disquieting, showing that reanalyses using es-
211 sentially the same data, and models that have been intercompared over decades, have significant
212 qualitative disagreements on climate time scales. Differences in the reanalyses in the northern
213 hemisphere are not so large, and are generally agreed to be the result of a much greater data
214 density. They remain, nevertheless, significant, as evidenced in the discussion of analysis in-
215 crements over the Arctic by Cullather and Bosilovich (2012). *Evidently, considerations of data*
216 *density and types dominate the reanalyses, with the models being of secondary importance.*

217 For climate studies, another major concern is the failure of the reanalyses to satisfy basic
218 global conservation requirements. So for example, Table 1 shows the global imbalances on a

³We employ “physics” in its conventional meaning as encompassing all of dynamics and thermodynamics.

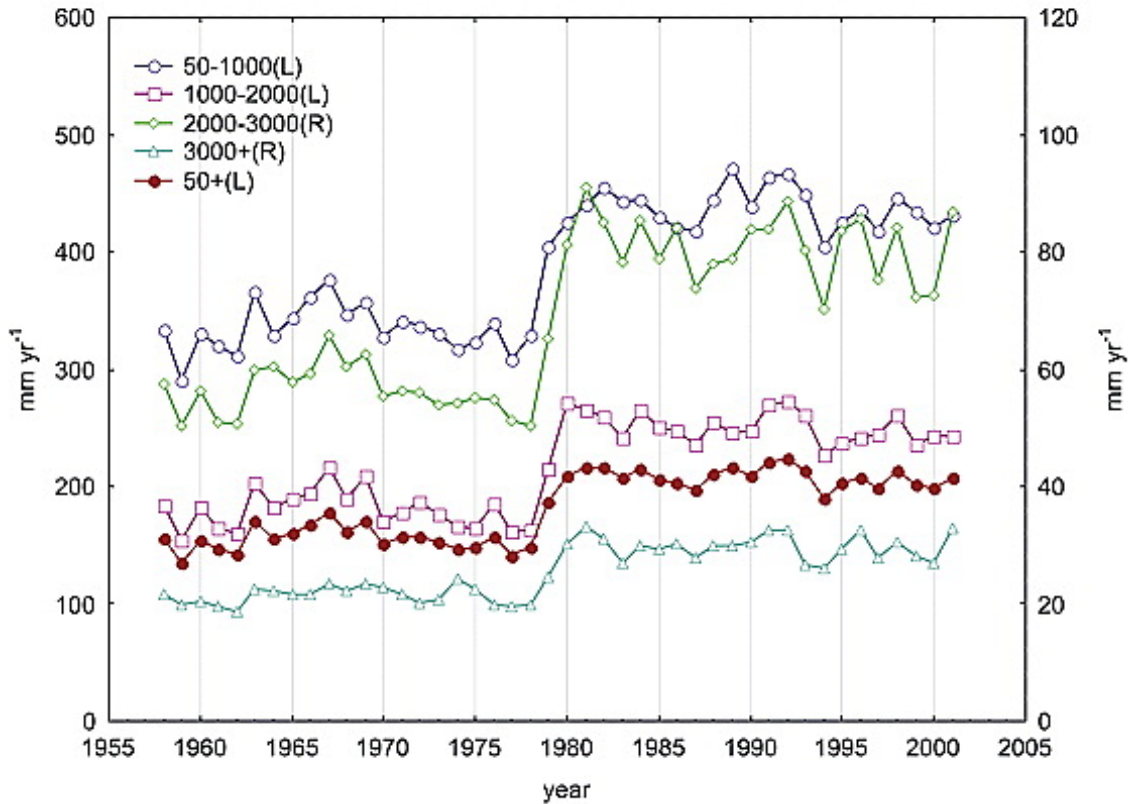


Figure 1: Mean annual precipitation minus evaporation over the Antarctic as a function of time in the ECMWF reanalysis ERA-40 showing the impact of new observations, in this case, the arrival of the polar orbiting satellites (from Bromwich et al., 2007). Different curves are for different elevations. The only simple inference is that the uncertainties must exceed the size of the rapid transition seen in the late 1970s. L and R identify whether the left or right axis is to be used for that curve.

{bromwich_etal

219 per year basis of several reanalysis products in apparent heating of the oceans and in the net
 220 freshwater flux from the atmosphere. Such imbalances can arise either because global constraints
 221 are not implied by the model equations, and/or because biased data have not been properly
 222 handled, or most likely, some combination of these effects is present. Trenberth and Solomon
 223 (1994) for example, note that the NCEP/NCAR reanalysis implies a meridional heat transport
 224 within continental land masses. “User beware” is the best advice we can give.

225 State estimation as defined in the ECCO context is a much more robust and tractable prob-
 226 lem than is, for example, prediction of future climate states. As is well known even to beginning
 227 scientists, extrapolation of very simple models can be extremely unstable, with interpolation ⁴,

⁴The commonplace term “interpolation,” is used in numerical analysis to imply that fitted curves pass exactly through data points—an inappropriate requirement here.

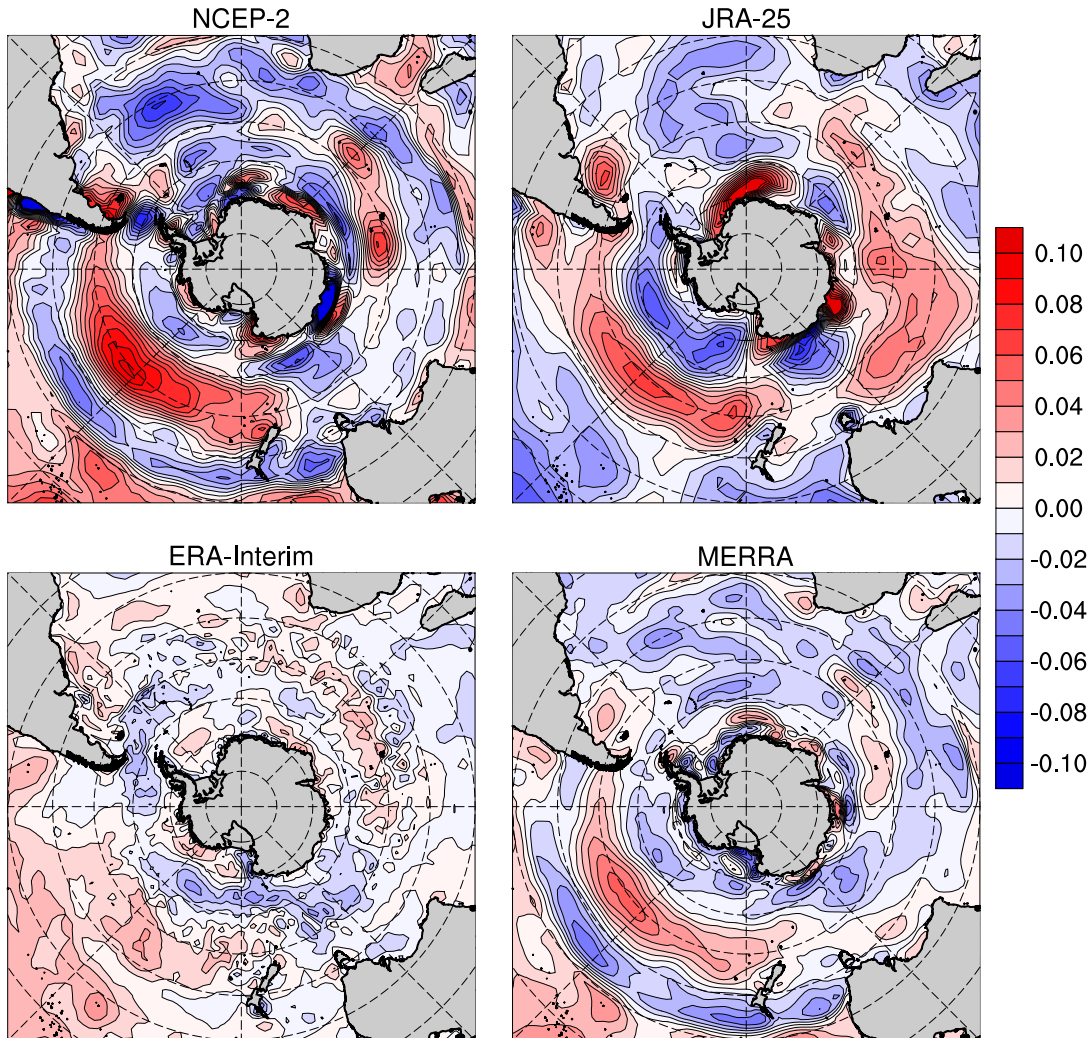


Figure 2: Calculated trends (meters/second/year) in the 10meter zonal wind fields at high southern latitudes from four different atmospheric reanalyses (D. Bromwich and J. P. Nicolas, of Ohio State University, private communication, 2010). Note particularly the different patterns in the Indian Ocean and the generally discrepant amplitudes. Because of the commonality of data sets, forecast models, and methodologies, the differences here must be lower bounds on the true uncertainties of trends. See Bromwich et al. (2011) for a description of the four different estimates.

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228 or curve-fitting, remaining robust. (A classical example is the use of a cubic polynomial to fit
 229 some noisy data, and which can be very effective. But one is advised *never* to use such a fit to
 230 extrapolate the curve; see Fig. 3). The ECCO process is effectively a temporal curve-fit of the
 231 WOCE-era data sets by a model and which, with some care to avoid data blunders, produces
 232 a robust result. It is the interpolating (“smoothing”) character, coupled with the expectation

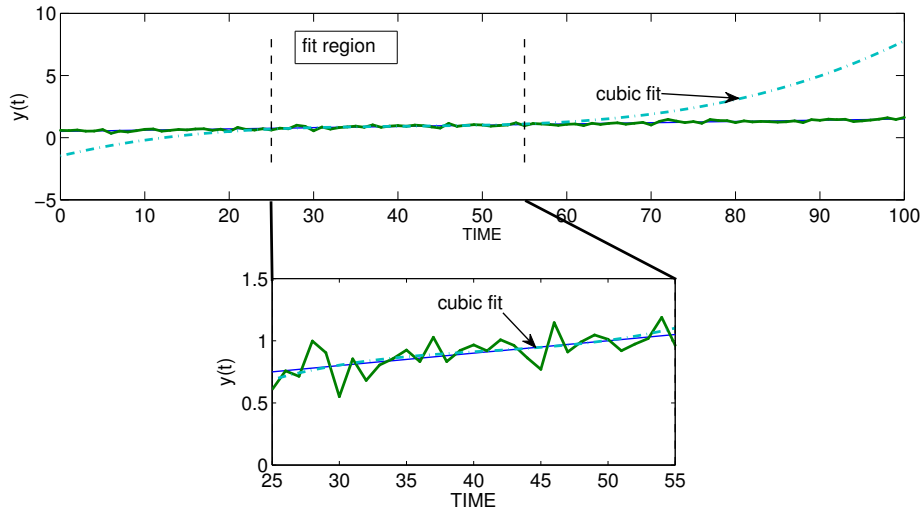


Figure 3: A textbook example of the robust interpolation of noisy data by a cubic polynomial and its gross instability when used to extrapolate. This analogue is a very simplified representation of the problem of extrapolating a GCM state into unobserved time spans.

{cubic_fit.eps}

233 of thermal wind balance over most of the domain, that produces confidence in the basic system
 234 products. As is well-known, least-squares methods tend to generate meaningless structures in
 235 unconstrained parts of the domain. Some regions of spatial extrapolation do exist here, depend-
 236 ing upon the time-varying distribution of observations, and although they tend to be limited in
 237 both space and time, detailed values there should be regarded skeptically.

238 *Terminological Note.* The observational community has lost control of the word “data,”
 239 and which has come to be used, confusingly, for the output of models, rather than having any
 240 direct relationship to instrumental values. In the context of reanalyses and state estimates
 241 involving both measurements and computer codes, the word generally no longer conveys any
 242 information. For purposes of this essay, “data” always represents instrumental values of some
 243 sort, and anything coming out of a GCM is a “model-value” or “model-datum,” or similar label.
 244 We recognize that models are involved in all real observations, even in such familiar values
 245 as those coming from e.g., a mercury thermometer, in which a measured length is converted
 246 to a temperature. Most readers can recognize the qualitative difference between conventional
 247 observations and the output of a 100,000+ line computer code.

248 4 Ocean State Estimates

249 The remainder of this paper is primarily devoted to a summary description and discussion of
250 some results of the Estimating the Circulation and Climate of the Ocean (ECCO) groups which,
251 beginning with Stammer et al. (2002), were directed at decadal and longer state estimates
252 satisfying known equations and using as much of the WOCE-era-and-beyond data as possible.
253 No claim is made that these estimates are definitive, nor that the discussion is comprehensive.
254 A number of other, superficially similar, estimates exist (Carton et al. 2000; Martin et al.,
255 2007; Hurlburt et al., 2009), but these generally have had different goals e.g., a fast approxi-
256 mate estimate primarily of the upper ocean, or prediction of the mesoscales over ocean basins.
257 Some weather forecasting centers have undertaken “operational oceanography” products closely
258 resembling atmospheric weather forecasts. To our knowledge, however, the ECCO estimates
259 are today the only ones specifically directed at physically continuous, dynamically consistent,
260 top-to-bottom estimates from a comprehensive data set.

261 A number of review papers exist that attempt to compare different such solutions (e.g.,
262 Carton and Santorelli, 2008; Lee et al., 2010) as though they were equivalent. But as the above
263 discussion tries to make clear, estimates are not equally reliable for all purposes and comparisons
264 make no sense unless their individual purposes are well understood. Although one could compare
265 a crop-dusting airplane to a jet fighter, and both have their uses, few would regard that effort
266 as helpful except as a vehicle for discussion of the highly diverse applications of aero-physics.
267 Thus a numerical scheme directed primarily at mesoscale prediction, and using a model not
268 conserving energy, may well be a useful tool for forecasting the trajectory of the Gulf Stream
269 over a few weeks, but it would be unsuited to a discussion of global ocean heat transports—a
270 useful model of which is, in turn, unsuitable for mesoscale interests. These other applications
271 are discussed in this volume by Schiller et al. (2012).

272 Originally, ECCO was meant primarily to be a demonstration of the practicality of its
273 approach to finding the oceanic state. When the first ECCO estimates did become available
274 (Stammer et al., 2002) they proved sufficiently useful even with that short duration and coarse
275 resolution, that a decision was made to continue with a gradually improving data set and
276 computer power. This review summarizes mainly what has been published thus far, but as
277 optimization is an asymptotic process, the reader should be aware that newer, and likely better,
278 solutions are being prepared continuously and the specific results here will have been refined in
279 the intervals between writing, publishing, and reading.

280 **4.1 Basic Notions**

281 As described above, most state estimation problems in practice are generically those of con-
 282 strained least-squares, in which one seeks to minimize objective or cost or misfit functions
 283 similar to Eq. (6) subject to the solution (including both the estimated state $\mathbf{x}(t)$, and the con-
 284 trols, $\mathbf{u}(t)$) of the model-time stepping equations.⁵ One approach, among many, to solving such
 285 problems is the method of Lagrange multipliers (MLM) dating back 200 years. This method is
 286 discussed at length in Wunsch (2006) and the references there. In a very brief summary, one
 287 “adjoins” the model equations using vectors of Lagrange multipliers, $\boldsymbol{\mu}(t)$, to produce a new
 288 objective function,

$$\begin{aligned}
 J' = & \sum_{m=0}^M (\mathbf{y}(t) - \mathbf{E}(t) \mathbf{x}(t))^T \mathbf{R}(t)^{-1} (\mathbf{y}(t) - \mathbf{E}(t) \mathbf{x}(t)) & (9) \quad \{\text{Jp}\} \\
 & + \sum_{m=0}^{M-1} \mathbf{u}(t)^T \mathbf{Q}(t)^{-1} \mathbf{u}(t) \\
 & - 2 \sum_{m=1}^{M-1} \boldsymbol{\mu}(t)^T [\mathbf{x}(t) - \mathbf{L}\mathbf{x}(t - \Delta t), \mathbf{B}\mathbf{q}(t - \Delta t), \boldsymbol{\Gamma}\mathbf{u}(t - \Delta t)], \\
 & t = m\Delta t, \quad m = 0, \dots, M
 \end{aligned}$$

289 Textbooks explain that the problem can now be treated as a conventional, unconstrained
 290 least-squares problem in which the $\boldsymbol{\mu}(t)$ are part of the solution. In principle, one simply does
 291 vector differentiation with respect to all of $\mathbf{x}(t)$, $\mathbf{u}(t)$, $\boldsymbol{\mu}(t)$, sets the results to zero, and solves
 292 the resulting “normal equations” (they are written out in Wunsch, 2006). J and J' are very
 293 general, and one easily adds e.g., internal model parameters such as mixing coefficients, water
 294 depths, etc. as further parameters to be calculated, thus rendering the problem one of combined
 295 state and parameter estimation.

296 *The entire problem of state estimation thus reduces to finding the stationary values of J' .*
 297 The large literature on what is commonly called the “adjoint method” (“4DVAR” in weather
 298 forecasting, where it is used only incrementally over short time-spans) reduces to coping with
 299 a very large set of simultaneous equations (and some are nonlinear). But as an even larger
 300 literature deals with solving linear and nonlinear simultaneous equations by many methods,
 301 ranging from direct solution, to downhill search, to Monte Carlo, etc., most of the discussion

⁵Advantages exist to using norms other than L_2 including those such as one and infinity norms commonly regarded as robust. These norms are not normally used in ocean and atmosphere state estimation or data assimilation systems because software development for parallel computers permitting computation at super-large dimensions has not yet occurred.

302 of adjoint methods reduces to technical details, many of which are complex, but which are
303 primarily of interest to computer-code constructors (Heimbach et al., 2005). Within the normal
304 equations, the time-evolution of the Lagrange multipliers is readily shown to satisfy a set of
305 equations usually known as the “adjoint” or “dual” model. This dual model can be manipulated
306 into a form having time run “backwards,” although that interpretation is unnecessary; see the
307 references.

308 One very interesting complication is worth noting: the description in the last two paragraphs
309 assumes one can actually differentiate J, J' . In oceanographic practice, that implies differentiat-
310 ing the computer code which does everything. The “trick” that has made this method practical
311 for GCMs is so-called automatic (or algorithmic) differentiation (AD), in which a software tool
312 can be used to produce the partial derivatives and their transposed values—in the form of an-
313 other software code (see, e.g., Giering and Kaminski, 1998; Griewank and Walther, 2008; Utke
314 et al., 2008). This somewhat bland statement hides a complex set of practical issues; see e.g.,
315 Heimbach et al. (2005) for discussion in the context of the MIT general circulation model (MIT-
316 gcm). Most of the difficult problems are of no particular concern to someone mainly interested
317 in the results.⁶

318 As discussed in more detail by Wunsch and Heimbach (2007), the central ECCO estimates
319 are based upon this Lagrange multiplier method, with the state estimates obtained from the
320 adjusted, but then *freely running*, MITgcm, as is required in our definition of state estimation.
321 At the time of this writing, most of the estimates have restricted the control variables (the ad-
322 justable parameters) to the initial conditions and the meteorological forcing, although following
323 exploratory studies by Ferreira et al. (2005), Stammer (2005) and Liu et al. (2012a), state
324 estimates are beginning to become public that also adjust internal model parameters, such as
325 isopycnal, thickness or vertical diffusion.

326 A full modern oceanic general circulation model (GCM or OGCM) such as that of Marshall
327 et al. (1997) as modified over subsequent years (e.g., Adcroft et al., 2004; Campin et al.,
328 2004), is a complex machine consisting of hundreds thousands of lines of code encompassing
329 the Navier-Stokes equations, the relevant thermodynamics, sea ice and mixed-layer sub-codes,
330 various schemes to represent motions below the model resolution (whatever it may be), and
331 further subsidiary codes for overflow entrainment, etc. Understanding such a model is a difficult
332 proposition, in part because different elements were written by different people over many years,
333 sometimes without full understanding of the potential interactions of the existing or future

⁶The situation is little different than that in ordinary ocean GCM studies. Technical details of time-stepping, storage versus recomputation, re-starts, etc. are very important and sometimes very difficult, but not often of consequence to most readers, except where the author necessarily calls attention to them.

334 subcomponents. Furthermore, various studies have shown the inevitability of coding errors (e.g.,
335 Basili et al., 1992) and unlike the situation with the real ocean, one is faced with determining
336 if some interesting or unusual behavior is real or an artifact of interacting, possibly very subtle,
337 errors. (Nature presumably never solves the incorrect equations; but observational systems do
338 have their own mysteries that must be understood: recent examples include the discovery of
339 systematic errors in fall rates to infer the depth of XBT data e.g., by Wijffels et al. (2008), and
340 calibration errors of pressure sensors onboard some of the Argo floats (Barker et al., 2011)).

341 By recognizing that most algorithms can be regarded as directed at the approximate solution
342 of a least-squares problem, one can exploit the two-hundred-year history of methodologies that
343 have emerged (e.g., Björck, 1996), substituting differing numerical algorithms where necessary.
344 For example Köhl and Willebrand (2002) and Lea et al. (2002) suggested that the Lagrange
345 multiplier method would fail when applied at high resolution to oceanic systems that had become
346 chaotic. Although such behavior has been avoided in oceanographic practice (Gebbie et al.,
347 2006; Hoteit et al., 2006; Mazloff et al. 2010), one needs to separate the *possible* failure of
348 a particular numerical algorithm to find a constrained minimum from the inference that no
349 minimum exists. If local gradient descent methods are not feasible in truly chaotic systems, one
350 can fall back on variations of Monte Carlo or other more global methods. Obvious failure of
351 search methods using local derivatives has had limited importance in oceanographic practice.
352 This immunity is likely a consequence of the observed finite time interval in the state estimation
353 problem, in which structures such as bifurcations are tracked adequately by the formally future
354 data, providing adequate estimates of the algorithmic descent directions. Systematic failure to
355 achieve an acceptable fit to the observations can lead to accepting the hypothesis that the model
356 should be rejected as an adequate representation. Potential model falsification is part of the
357 estimation problem, and is the pathway to model improvement.

358 Modern physical oceanography is largely based upon inferences from the thermal wind, or
359 geostrophic-hydrostatic, equations. Scale analyses of the primitive equations (e.g., Pedlosky,
360 1987; Vallis, 2006; Huang, 2010) all demonstrate that apart from some very exceptional regions
361 of small area and volume, deviations from geostrophic balance are slight. This feature is simul-
362 taneously an advantage and a liability. It is an advantage because any model, be it analytical
363 or numerical must, to a first approximation satisfy the linear thermal wind equations. It is a
364 liability because it is only the deviations which define the governing physics of the flow main-
365 tenance and evolution, and which are both difficult to observe and to compute with adequate
366 accuracy. In the present context, one anticipates that over the majority of the oceanic volume,
367 any plausible model fit to the data sets must be, to a good approximation, a rendering of the
368 ocean circulation in geostrophic, hydrostatic, balance, with Ekman forcing, and volume or mass

369 conservation imposed regionally and globally as an automatic consequence of the model config-
370 uration. The most visible ageostrophic physics are the variability, seen as slow, accumulating,
371 deviations from an initial state.

372 4.2 The Observations

373 Data sets used for many (not all) of the ECCO family of solutions are displayed in Table 2. As
374 noted in the Introduction, they are of very diverse type, geographical and temporal distribution,
375 and with very different accuracies and precisions.

376 As is true of any least-squares solution, no matter how it is obtained, the results are directly
377 dependent upon the weights or error variances assigned to the data sets. An over-estimate of
378 the error corresponds to the suppression of useful information; an under-estimate to imposing
379 erroneous values and structures. Although an unglamorous and not well-rewarded activity, a
380 quantitative description of the errors is essential and is often where oceanographic expertise is
381 most central. Partial discussions are provided by Stammer et al. (2007), Ponte et al. (2007),
382 Forget and Wunsch (2007), and Ablain et al. (2009). Little is known about the space-time
383 covariances of these errors, information, which if it were known, could improve the solutions
384 (see Weaver and Courtier, 2001, for a useful direction now being used in representing spatial
385 covariances). Model errors, which dictate how well estimates should fit to hypothetical perfect
386 data, are extremely poorly known and are generally added to the true data error—as in linear
387 problems the two types of error are algebraically indistinguishable.

388 5 Global Scale Solutions

389 Solutions of this type were first described by Stammer et al. (2001, 2002, 2003) and were
390 computed on a $2^\circ \times 2^\circ$ grid with 22 vertical levels. As the computing power increased, a shift
391 was made to a $1^\circ \times 1^\circ$, 23-level solution and that, until very recently, has remained the central
392 vehicle for the global ECCO calculations. Although some discrepancies continue to exist in the
393 ability to fit certain data types, these solutions (Wunsch and Heimbach, 2007) based as they are
394 on geostrophic, hydrostatic balance over most of the domain, were and are judged adequate for
395 the calculation of large-scale transport and variability properties. The limited resolution does
396 mean that systematic misfits were expected, and are observed, in special regions such as the
397 western boundary currents. Often the assumed error structures of the data are themselves of
398 doubtful accuracy.

399 As noted above, Ganachaud (2003a) inferred that the dominant error in trans-oceanic trans-
400 port calculations of properties arose from the temporal variability. Perhaps the most important

401 lesson of the past decade has been the growing recognition of the extent to which temporal
402 aliasing is a serious problem in calculating the oceanic state. For example, Figs. 4, 5 display the
403 global meridional heat and fresh water transport as a function of latitude along with their stan-
404 dard errors computed from the monthly fluctuations. The figures suggest that errors inferred
405 from hydrography are under-estimated (and error estimates of the non-eddy resolving ECCO
406 estimates are themselves lower bounds of the noise encountered in the real ocean). The classical
407 oceanographic notion that semi-synoptic sections are accurate renderings of the time-average
408 properties, while having some qualitative utility, has now to be painfully abandoned—an es-
409 sential step if the subject is to be a quantitative one. Temporal effects are most conspicuous
410 at low latitudes, but in many ways, the difficulty is greatest at high latitudes: the long time
411 scales governing behavior there mean that the hydrographic structure is very slowly changing,
412 requiring far longer times to produce an accurate time-mean. In other words, a 10-year average
413 at 10°N will be a more accurate estimate of the longer-term mean than one at 50°N . Even this
414 comment begs the question of whether a stable long-term mean exists, or whether the system
415 drifts over hundreds and thousands of years? This latter is a question concerning the frequency
416 spectrum of oceanic variability, and which is very poorly known at periods beyond a few years.

417 For the 19+ years now available in the global state estimates, most of the large-scale prop-
418 erties, including the time variations, are stable from one particular set of assumptions to others,
419 probably as a consequence of the dominance of overall geostrophic balance and the comparatively
420 well-sampled hydrography and altimetric slopes. They are thus worth analyzing in detail. The
421 intricacies of the global, time-varying ocean circulation are a serious challenge to the summariz-
422 ing capabilities of authors. A full discussion, however, of the global state estimates becomes a
423 discussion of the complete three-dimensional time-varying ocean circulation, a subject requiring
424 a book, if not an entire library, encompassing distinctions amongst time and space scales, geo-
425 graphical position, depth, season, trends, the forcing functions (controls). No such synthesis is
426 attempted here! Instead we can only give a bit of the flavor of what can and has been done with
427 the estimates along with enough references for the reader to penetrate the wider literature.

428 Note too, as discussed e.g., by Heimbach et al. (2011) and other, earlier efforts, the Lagrange
429 multipliers are the solution to the dual model. As such, they are complete solutions in three
430 spatial dimensions and time, and convey the sensitivity of the forward model to essentially any
431 parameter or boundary or initial condition in the system. The information content of the dual
432 solutions is very large—representing not only the sensitivities of the solution to the data and
433 model parameters and boundary and initial conditions, but also the flow of information through
434 the system. Analyzing the dual solution does, however, require the same three-dimensional time-
435 dependent representations of any full GCM, and these elements of the state estimates remain

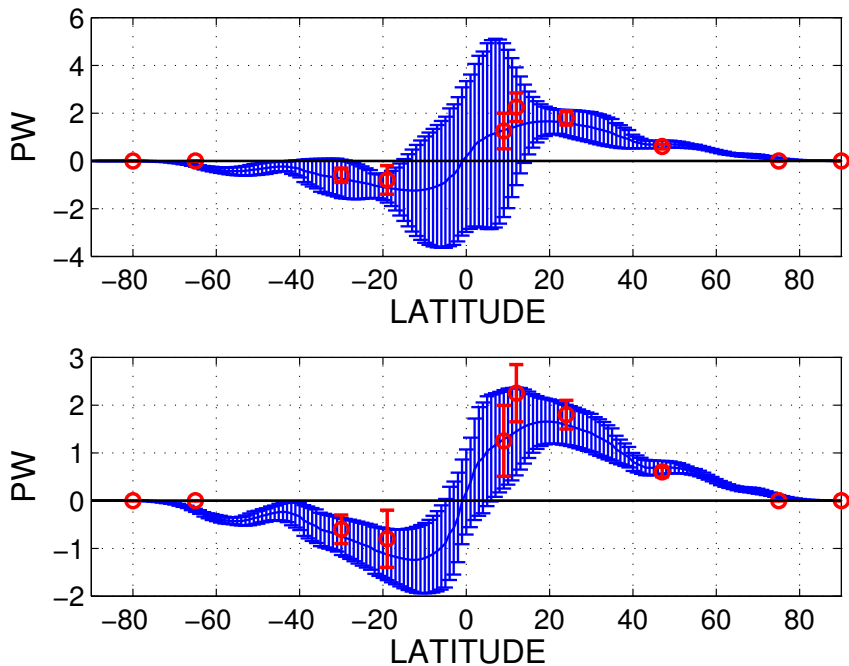


Figure 4: Global meridional heat transport in the ocean from ECCO-Production version 4 (G. Forget, private communication, 2011). Upper panel shows the standard error including the annual cycle and the lower one, with the annual cycle removed—as being largely predictable. Possible systematic errors are not included. Red dots with error bars are estimates from Ganachaud and Wunsch (2002). Note that the WOCE-era hydrographic survey failed to capture the southern hemisphere extreme near 10°S, thus giving an exaggerated picture of the oceanic heat transport asymmetry about the equator.

{heat_transpor

436 greatly under-exploited at the present time.

437 5.1 Summary of Major, Large-Scale Results

438 None of the results obtained so far can be regarded as the final state estimate: obtaining fully
 439 consistent misfits by the model to the observations has never been achieved (see the residual
 440 misfit figures in the references). Misfits linger for a variety of reasons, including the sometimes
 441 premature termination of the descent algorithms before full optimization, mis-representation of
 442 the true model or data errors, or selection of a local rather than a global minimum in the major
 443 nonlinear components of the model. As with all very large nonlinear optimization problems,
 444 approach to the “best” solution is asymptotic. With these caveats, we describe some of the more
 445 salient oceanographic features of the recent solutions, with no claim to being comprehensive.
 446 Note that results from a variety of ECCO-family estimates are used, largely dictated by the

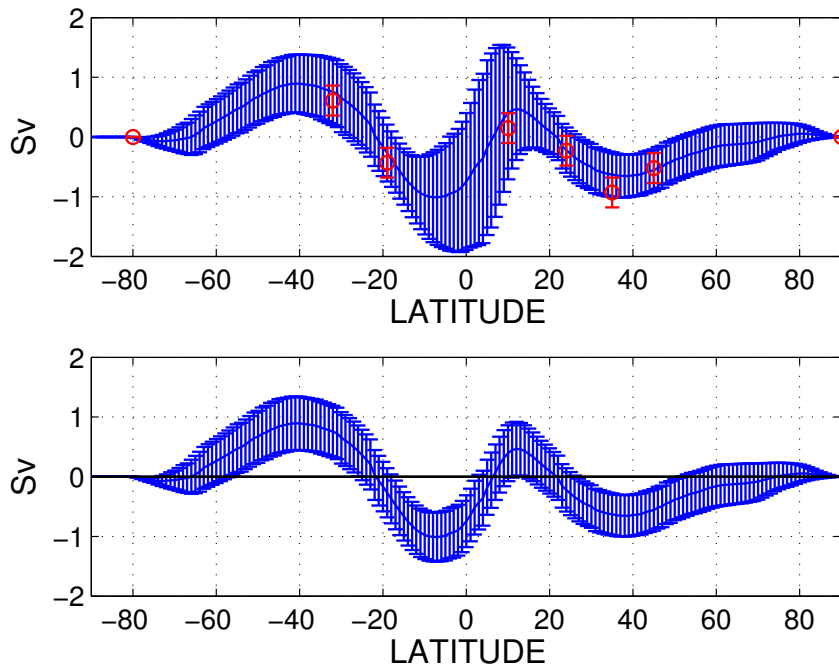


Figure 5: Same as Fig. 4 except for the freshwater transport (G. Forget, private communication, 2011). Upper panel shows standard errors that include the seasonal cycle, and the lower without the seasonal cycle. Red dots are again from Ganachaud and Wunsch (2002).

{fw_transport_

447 particular problem that was the focus of the calculation.

448 *Volume, Enthalpy, Freshwater Transports and Their Variability*

449 The most basic elements describing the ocean circulation and its large-scale variability are
 450 usually the mass (or volume, which is nearly identical) transports. Stammer et al. (2001,
 451 2002, 2003) depicted the basic global-scale elements of the mass transport as averaged over
 452 the duration of their estimates. A longer duration estimate (v3.73) has been used (Fig. 6) to
 453 compute the vertically integrated volume stream function. We reiterate that diagrams such as
 454 this one are finite duration averages whose relationship to hypothetical hundred year or longer
 455 climatologies remains uncertain.

456 Fig. 7 shows the zonally integrated and vertically accumulated meridional transport as a
 457 function of depth and ocean. The very large degree of temporal variability can be seen in Fig. 4
 458 from a new fully global solution which is about to become available online at the time of writing
 459 (ECCO-Production version 4; see Table 3) with error bars derived from the temporal variances.
 460 These time averages have been an historically important goal of physical oceanography, albeit
 461 estimates derived from unaveraged data were commonly assumed without basis to accurately

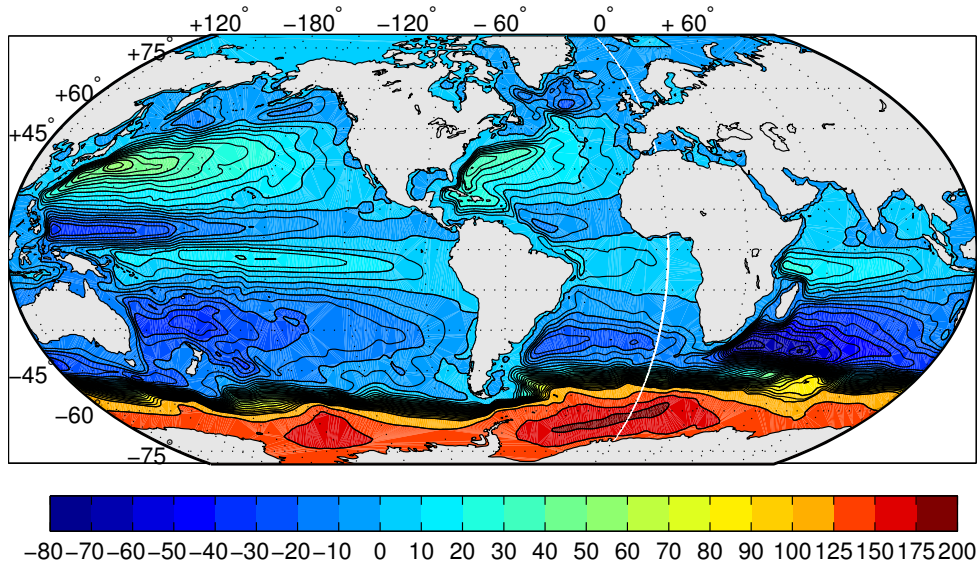


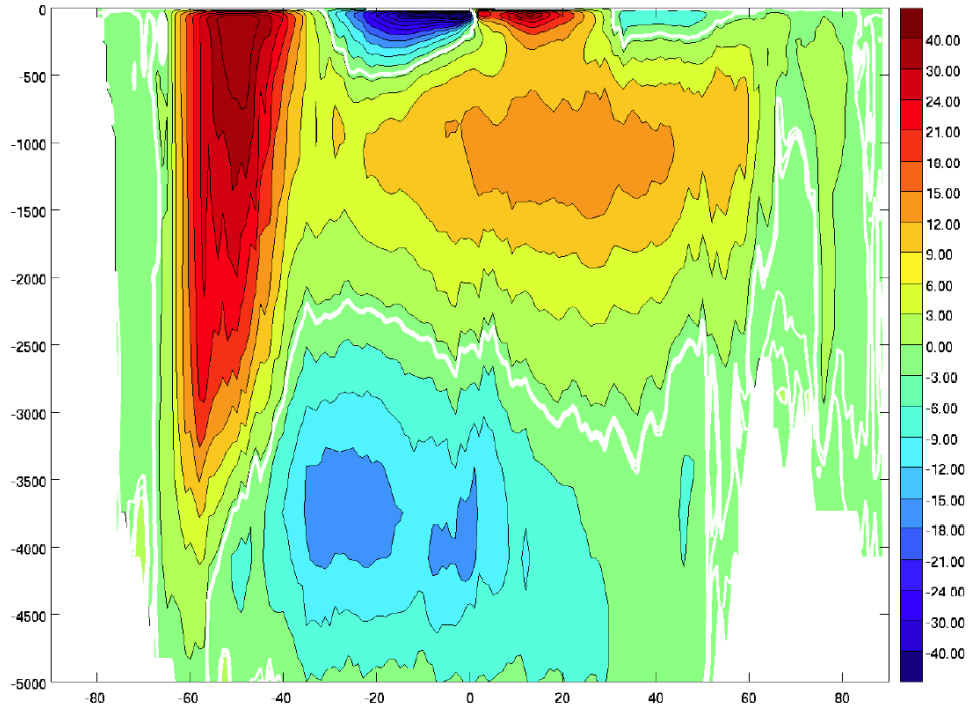
Figure 6: The top-to- bottom transport stream function from ECCO v3.73 (Wunsch, 2011). Qualitatively, the wind-driven gyres dominate the result, with the intense transports in the Southern Ocean particularly conspicuous.

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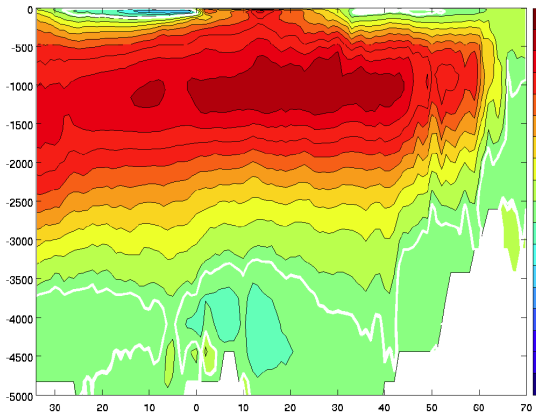
462 depict the true time-average. Perhaps the most important utility of the existing state estimates
 463 has been the ability, at last, to estimate the extent of the time-variability taking place in the
 464 oceans (Wunsch and Heimbach, 2007, 2009, 2012). Withheld, direct in situ observations in
 465 a few isolated regions (Kanzow et al., 2009; Baehr, 2010) are consistent with the inference
 466 that even volume transports integrated across entire ocean basins have a large and qualitative
 467 temporal variability. More generally, mooring data and the now almost 20-year high resolution
 468 high accuracy altimetric records all show the intense variations that exist everywhere. With
 469 ECCO-like systems, syntheses of these data sets are now possible.

470 *The Annual Cycle*

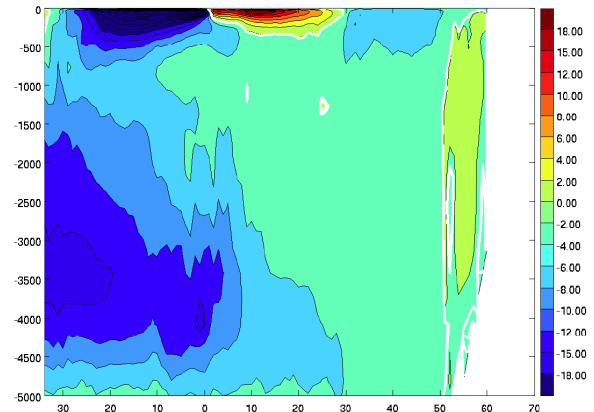
471 The annual cycle of oceanic response is of interest in part because the ultimate forcing
 472 function (movement of the sun through the year) is very large and with very accurately known
 473 structure. In practice, that forcing is mediated through the very complex atmospheric annual
 474 changes, and understanding how and why the ocean shifts seasonally on a global basis is a
 475 difficult problem. Using the ECCO state estimates, Vinogradov et al. (2008) mapped the
 476 amplitude and relative contributions for salt and heat of the annual cycle in sea level (Figs.



(a) Global



(b) Atlantic Ocean



(c) Pacific & Indian Ocean

Figure 7: Mean (1992-2010) of the meridional volume transport stream function in Sverdrups ($\text{Sv} = 10^6 \text{m}^3/\text{s}$) from ECCO-Production version 4 (Wunsch and Heimbach, 2012; Forget et al., in prep. 2012). Panel (a) is the global result; panels (b,c) are the Atlantic, and the combined Indo-Pacific, respectively. Note the complex equatorial structure, and that this representation integrates out a myriad of radically different dynamical sub-regimes. In the Southern Ocean, interpretation of zonally integrated Eulerian means requires particular care owing to the complex topography and relatively important eddy transport field.

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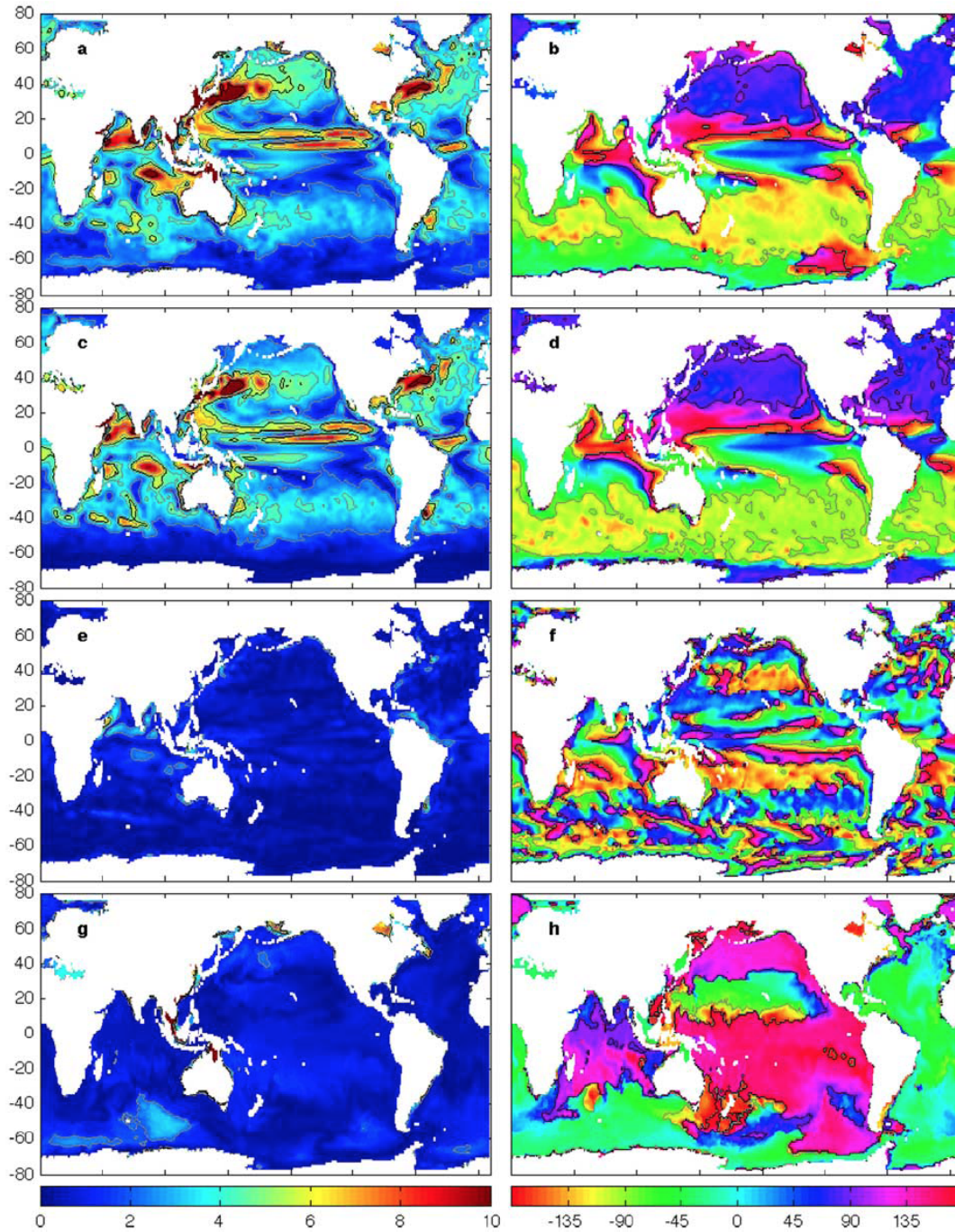


Figure 8: From Vinogradov et al. (2008) showing the annual cycle in sea level from ECCO Climate State v2.177. Left column is the amplitude in cms and the right column the phase. From top-to-bottom, they are the surface elevation (a,b), the thermosteric component (c,d), the halosteric component (e,f), and at bottom, the bottom pressure (g,h).

{vinogradov_et

477 8. The importance of the annual cycle, more generally, is visible in Fig. 4, 5 as the large
478 contribution to the standard errors.

479 *Sea Level Change*

480 The sea surface height is simultaneously a boundary condition on the oceanic general circula-
481 tion and a consequence of that circulation. Because of the intense interest in possible large-scale
482 changes in its height and the potential shifts in vulnerability to storm surges, and associated
483 issues such as ecosystem and freshwater reservoir declines, the ECCO state estimation system
484 has been used to estimate the shifts taking place in the era since 1992 (Wunsch et al., 2007).
485 A summary of a complex subject is that sea level change is dominated by regional variations
486 more than an order of magnitude larger than the putative global average, and arising primarily
487 from wind field shifts. Varying spatial contributions from competing exchanges of freshwater
488 and heat with the atmosphere and the extremely inhomogeneous (space and time) in situ data
489 sets render the global mean and its underlying causes far more uncertain than some authors
490 have claimed.

491 At the levels of accuracy appearing to be required, very careful attention must now be paid
492 to modeling issues such as water self-attraction and load (Vinogradova et al., 2011, Kuhlmann et
493 al., 2011) not normally accounted for in OGCMs. Conventional approximations to the moving
494 free-surface boundary conditions generate systematic errors no longer tolerable (e.g., Huang,
495 1993; Wunsch et al., 2007). Usefully accurate sea level estimation over multiple decades may
496 be the most demanding requirement on both models and data sets now facing oceanographers
497 (Griffies and Greatbatch, 2012). The global means are claimed by some to have accuracies
498 approaching a few tenths of a millimeter per year—an historically extraordinary requirement on
499 any ocean estimate. Despite widely publicized claims to the contrary (e.g., Cazenave and Remy,
500 2011; Church et al., 2011), state estimate results suggest that at the present time, the global
501 observing system appears to be insufficient to provide robust partitioning amongst heat content
502 changes, land and ice sheet runoff, and large-scale shifts in circulation patterns. A particular
503 difficulty pertains to the deep ocean, below depths measured by the Argo array, where the
504 distinction between apparent changes occurring (Kouketsu, 2011) and the significant deep eddy
505 variability (Ponte, 2012) remains obscure due to poor observational coverage. Claims for closed
506 budget elements involve accuracies much coarser than are stated for the total value. ⁷

507 *Biogeochemical Balances*

⁷We have omitted here the distinction between absolute sea level with respect to the geoid, and relative sea level measured by tide gauges, and ignored processes associated with the unloading of the solid Earth from ice sheet shrinkage. None of these is represented in current ocean or climate models (e.g., Munk, 2002; Milne et al., 2009).

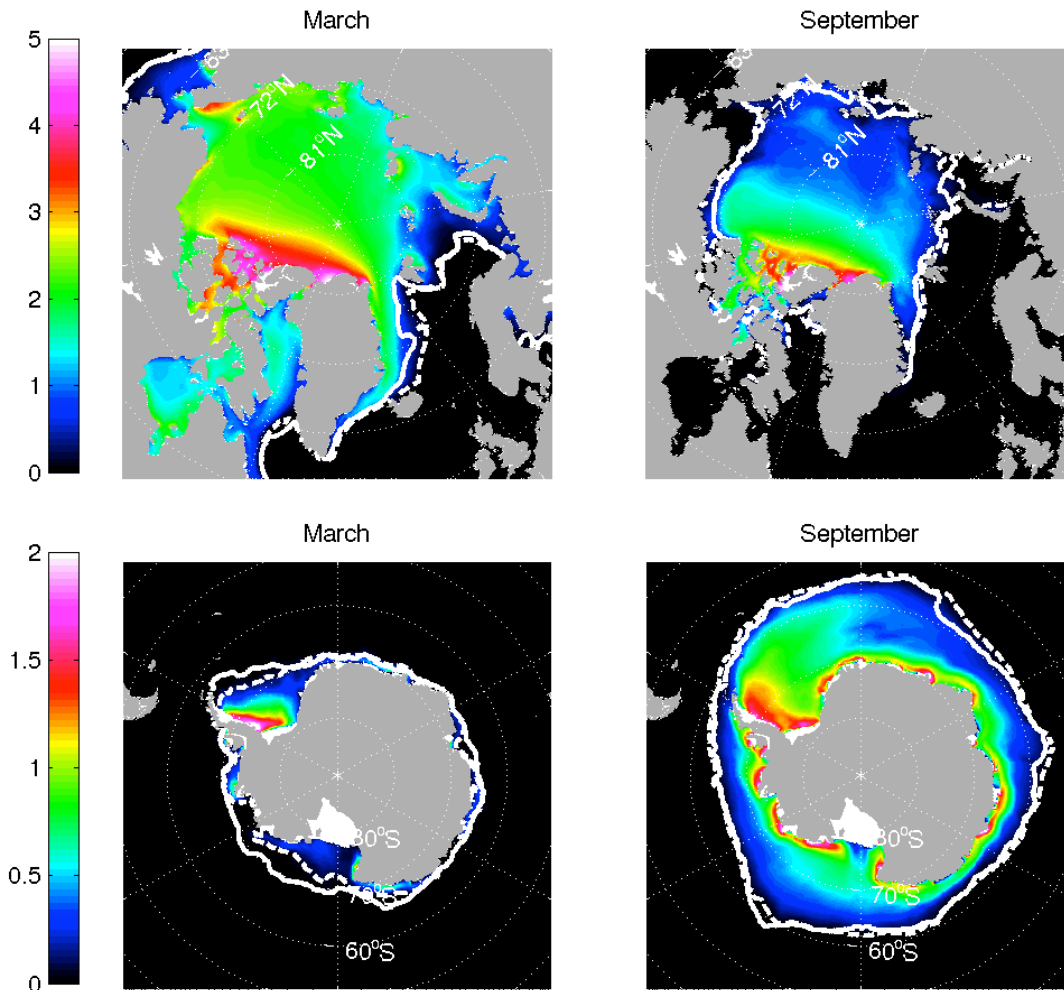


Figure 9: 1992-2002 mean March (left) and September (right) effective ice thickness distributions (in meters) for northern (top) and Southern (bottom) hemispheres. Obtained from a global eddy-permitting ECCO2 simulation, for which a set of global parameters has been adjusted. Also indicated are the ice edge (15% ice concentration isoline) inferred from the model (dashed line) and from satellite-retrieved passive microwave radiometry (solid line). From Losch et al. (2010).

{cube76marsepi

508 From the adjoint of the tracer concentration sub-model of the ECCO system Dutkiewicz et
 509 al. (2006) calculated the sensitivity of the nutrient production in the system to iron enrichment.
 510 This work is representative of the use of dual solutions to probe large complex models in any
 511 field. They found a strong dependence upon the available light, and that the tropical ocean
 512 had the greatest sensitivity to iron limitation. Among other considerations, these inferences are
 513 important in the erstwhile debate over whether iron fertilization makes any sense for control of
 514 atmospheric CO₂.

515 Woloszyn et al. (2011) used the ECCO higher-resolution Southern Ocean State Estimates

516 (SOSE) of Mazloff et al. (2010) to demonstrate the great importance of adequate resolution
517 in calculating carbon exchange between the atmosphere and ocean. The same configuration
518 was adopted by Ito et al. (2010) to describe the Ekman layer contribution to the movement of
519 carbon dioxide.

520 The emerging field of microbial oceanography seeks a zeroth-order understanding of the
521 biogeography and diversity of marine microbes. Coupling between ocean physics and ecology is
522 being explored through the use of ECCO state estimates which drive models of marine ecology
523 (e.g., Follows et al., 2007; Follows and Dutkiewicz, 2011). Crucial requirements of the estimates
524 are (1) to be in sufficiently close agreement with the observed physical ocean state such as to
525 reduce uncertainties in the coupled models from the physical component, and (2) to furnish an
526 evolution of the physical state in agreement with conservation laws.

527 *Sea Ice*

528 The importance of sea ice to both the ocean circulation and climate more generally has
529 become much more conspicuous in recent years. Sea ice models have been developed within
530 the state estimation framework as fully coupled sub-systems influenced by and influencing the
531 ocean circulation (Menemenlis et al., 2005; Losch et al., 2010). By way of example, Fig. 11, taken
532 from Losch et al. (2010) depicts 1992-2002 mean March and September effective ice thickness
533 distributions representing the months of maximum and minimum ice cover in both hemispheres.
534 Also shown are the modeled and observed ice edge, represented as 15% isolines of the fractional
535 sea ice concentrations (0 to 100%). The results were obtained from an early version of the
536 ECCO2 eddy-permitting alternative optimization method using Green functions (Menemenlis
537 et al., 2005a,b) on the cubed-sphere grid at 18 km horizontal resolution (see Table 3). More
538 detailed studies focussing on the Arctic were carried out with similar and higher-resolution (4
539 km) configurations (Nguyen et al., 2011, 2012), but with a very limited control space available
540 for parameter adjustment via the Green function approach.

541 A comprehensive step toward full coupled ocean-sea ice estimation, in which both ocean and
542 sea ice observation were synthesized, was made by Fenty and Heimbach (2012a,b) for the limited
543 region of the Labrador Sea and Baffin Bay. Fig. 10a shows an annual cycle of total sea ice area
544 in the domain from observations, the state estimate, and the unconstrained model solution.
545 Also shown are the remaining misfits, as evidence of the random nature of the residuals, as
546 required by theory, Eqns. (2) and (3b). An important result of that study is the demonstration
547 that adjustment well within their prior uncertainties in the high-dimensional space of uncertain
548 surface atmospheric forcing, patterns can achieve an acceptable fit between model and observa-
549 tion, placing stringent requirements on process studies which aim to discriminate between model
550 errors and forcing deficiencies.

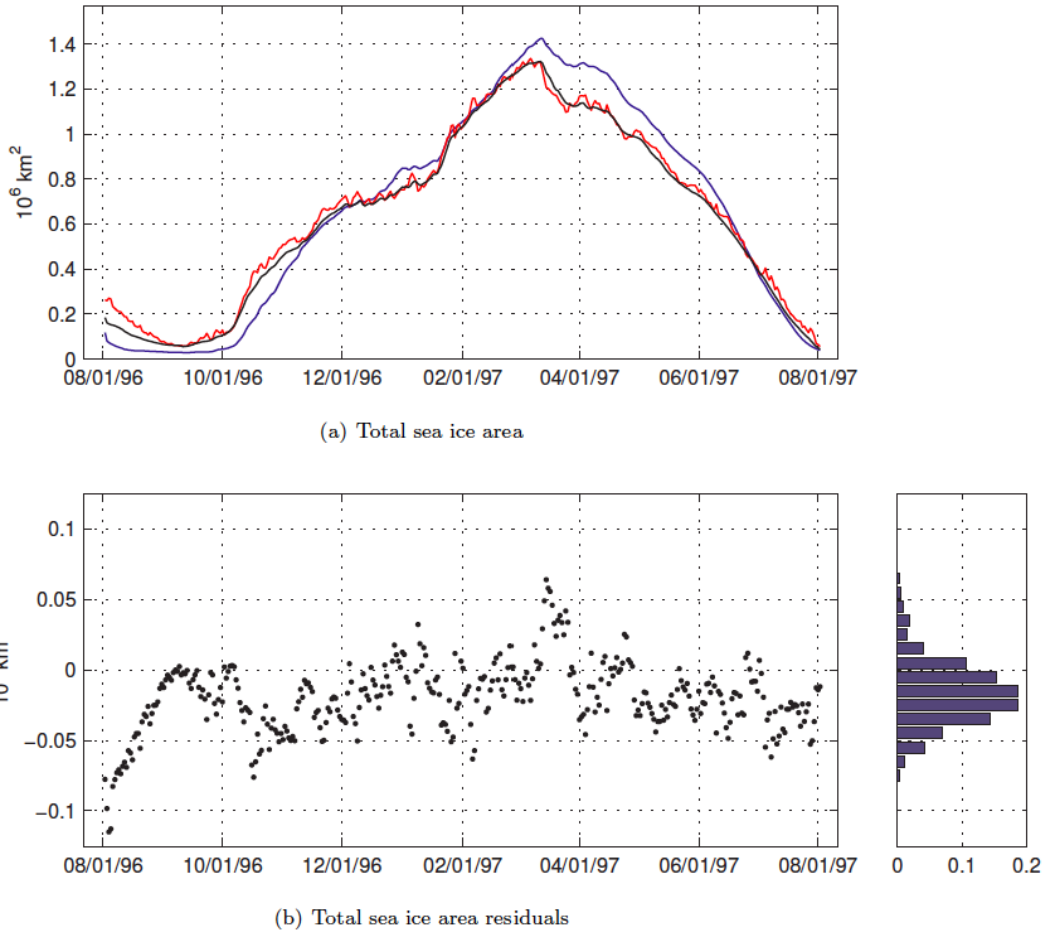


Figure 10: (Top): Annual cycle from August 1996 to July 1997 of daily-mean total sea ice area in the Labrador Sea and Baffin Bay from observations (red), a regional state estimate (black), and the unconstrained model solution (blue). (Bottom): Residual misfits between estimated and observed sea ice area and its frequency of occurrence histogram (right panel). Taken from Fenty et al. (2012a).

{fenty-labsea-

551 As in the discussion of biogeochemical balances above, the adjoint or dual solution of the
 552 coupled ocean-sea ice model can provide detailed sensitivity analyses. Heimbach et al. (2010)
 553 used the dual solution to study sensitivities of sea ice export through the Canadian Arctic
 554 Archipelago to changes in atmospheric forcing patterns in the domain. Kauker et al. (2009)
 555 investigated the causes of the 2007 September minimum in Arctic sea ice cover in terms of
 556 sensitivities to atmospheric forcing over the preceding months. A similar sensitivity study on
 557 longer time scales is shown in Fig. 11 of solid (sea ice and snow) freshwater export through Fram
 558 Strait for two study periods, January 1989 to September 1993, and January 2003 to September
 559 2007 (unpublished work). The objective function was chosen to be the annual sea ice export
 560 between October 1992 and September 1993, and October 2006 and September 2007. Export

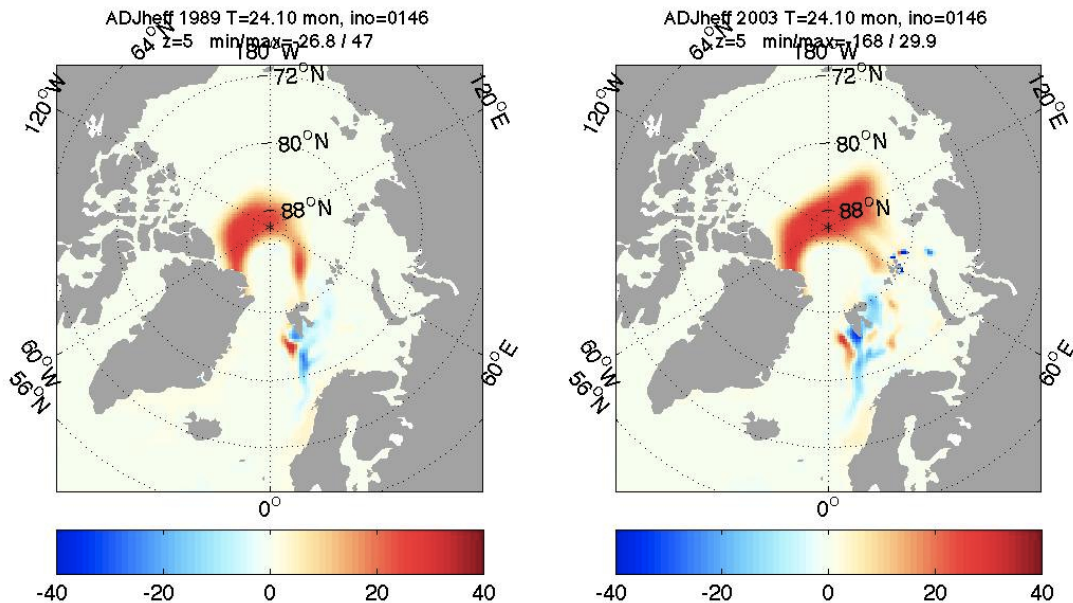


Figure 11: Sensitivity of sea ice export through Fram Strait to changes in effective sea ice thickness 24 months back in time. Two integration periods were considered, Jan. 1989 to Sep. 1993 (left) and Jan. 2003 to Sep. 2007 (right). The objective function is annual sea ice export between October 1992 and September 1993 (left), and October 2006 and September 2007 (right).

{fram_export_2

561 sensitivities to changes in effective sea ice thickness, 24 months prior to September 1993 and
 562 2007, respectively, are shown. The dominant patterns are positive sensitivities upstream of
 563 Fram Strait, and for which an increase in ice thickness will increase ice export at Fram Strait
 564 24 months later. (Spurious patterns south of Svalbard have been attributed to masking errors
 565 in the sea ice adjoint model and were corrected in Fenty and Heimbach, 2012a.) Sensitivities
 566 are linearized around their respective states, and depend on the state trajectory. The extended
 567 domain of influence for the 2007 case compared to 1993 suggests more swift transport conditions
 568 in the central Arctic, possibly due to favorable atmospheric conditions, or to weaker sea ice, or
 569 both.

570 *Ice Sheet-Ocean Interactions*

571 The intense interest in sea level change and the observed acceleration of outlet glaciers spilling
 572 into narrow deep fjords in Greenland and ice streams feeding vast ice shelves in Antarctica
 573 (e.g., Payne et al., 2004; Alley et al., 2005; Shepherd and Wingham, 2007; Pritchard et al.,
 574 2009; Rignot et al., 2011) has led to inferences that much of the ice response may be due
 575 to regional oceanic warming at the glacial grounding lines, an area termed by Munk (2011)
 576 “*this last piece of unknown ocean*”. One such region is the Amundsen Sea Embayment in West
 577 Antarctica (Fig. 12, taken from Schodlok et al., 2012), where the ocean is in contact with

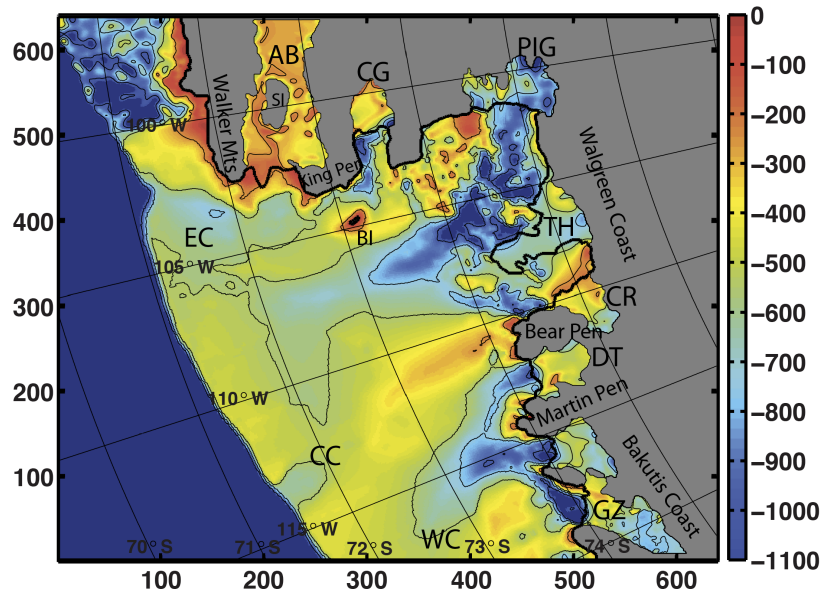


Figure 12: (From Schodlok et al., 2012): bottom topography (in meters) of the Amundsen Sea Embayment, West Antarctica, with thick black lines delineating the edge of several large ice shelves which buttress the following glaciers grounded deep below sea level: Abbot (AB), Cosgrove (CG), Pine Island Glacier (PIG), Thwaites (TH), Crosson (CR), Dotson (DT), and Getz (GZ). Also indicated are prominent topographic features, such as Sherman Island (SI), Burke Island (BI), Eastern Channel (EC), Central Channel (CC), and Western Channel (WC).

{pig_schodlok_

578 several large shelves, among which Pine Island Ice Shelf (PIIS) and Glacier (PIG) exhibits one
 579 of the largest changes in terms of ice sheet acceleration, thinning, and mass loss. Recent, and
 580 as yet incomplete model developments have been directed at determining the interactions of
 581 changing ocean temperatures and ice sheet response, and for the purpose of inclusion into the
 582 coupled state estimation system (Losch, 2008). Simulated melt rates under PIIS are depicted in
 583 Fig. 13, but cannot be easily measured directly. A first step toward their estimation in terms of
 584 measured hydrography has been undertaken by Heimbach and Losch (2012) who developed an
 585 adjoint model complementing the sub-ice shelf melt rate parameterization. By way of example,
 586 Fig. 14 depicts transient sensitivities of integrated melt rates (Fig. 13) to changes in ocean
 587 temperatures. The spatial inhomogeneous patterns have implications for the interpretation of
 588 isolated measurements and optimal observing design.

589 The critical dependence of sub-ice shelf cavity circulation and melt rates to details of the
 590 bathymetry and grounding line position noted by Schodlok et al. (2012) revives the issue of
 591 bottom topography as a dominant control on ocean circulation and the necessity for its inclusion
 592 into formal estimation systems (Losch and Wunsch, 2003; Losch and Heimbach, 2007).

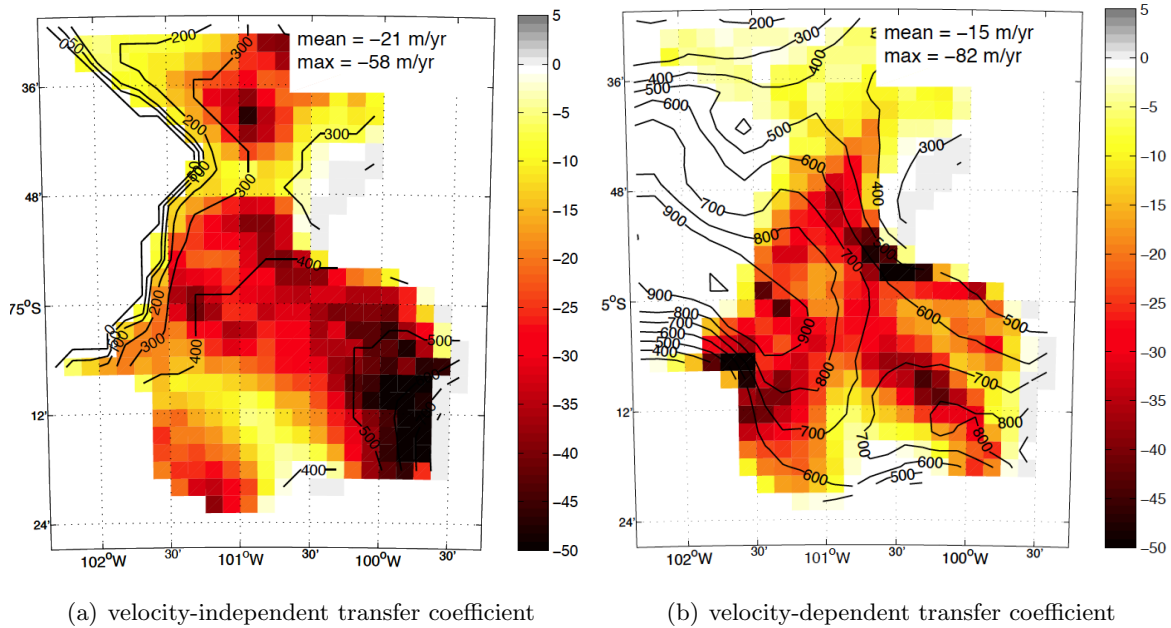


Figure 13: Simulated melt rates (colors, in meters/year) under Pine Island Ice Shelf (PIIS) derived from variants of the Holland and Jenkins (1999) melt rate parameterization, using either velocity-independent (a) or velocity-dependent transfer coefficients (from Dansereau, 2012). Large melt rates correspond to either locations deep inside the cavity where the ice shelf is in contact with the warmest Circumpolar Deep Waters, or to locations of highest flow at the ice shelf-ocean interface. Direct measurement of melt rates is challenging, making robust inferences difficult.

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593 *Air-Sea Transfers and Property Budgets*

594 By definition, state estimates permit calculations up to numerical accuracies of global budgets
 595 of energy, enthalpy, etc. Many of these budgets are of interest for the insight they provide into
 596 the forces powering the ocean circulation. Josey (2012, this volume) discusses estimates of the
 597 air-sea property transfers using the ECCO estimates. As an example, Fig. 15 is an estimate
 598 by Stammer et al. (2004) of the net air-sea transfers of fresh water. That paper compares this
 599 estimate to other more ad hoc calculations and evaluates its relative accuracy.

600 As examples of more specific studies using the state estimates, we note only Piecuch and
 601 Ponte (2011, 2012) who examined the role of transport fluctuations on the regional sea level and
 602 oceanic heat content distribution, and Roquet et al. (2011) who used them to depict the regions
 603 in which mechanical forcing by the atmosphere enters into the interior geostrophic circulation.
 604 Many more such studies are expected in the future.

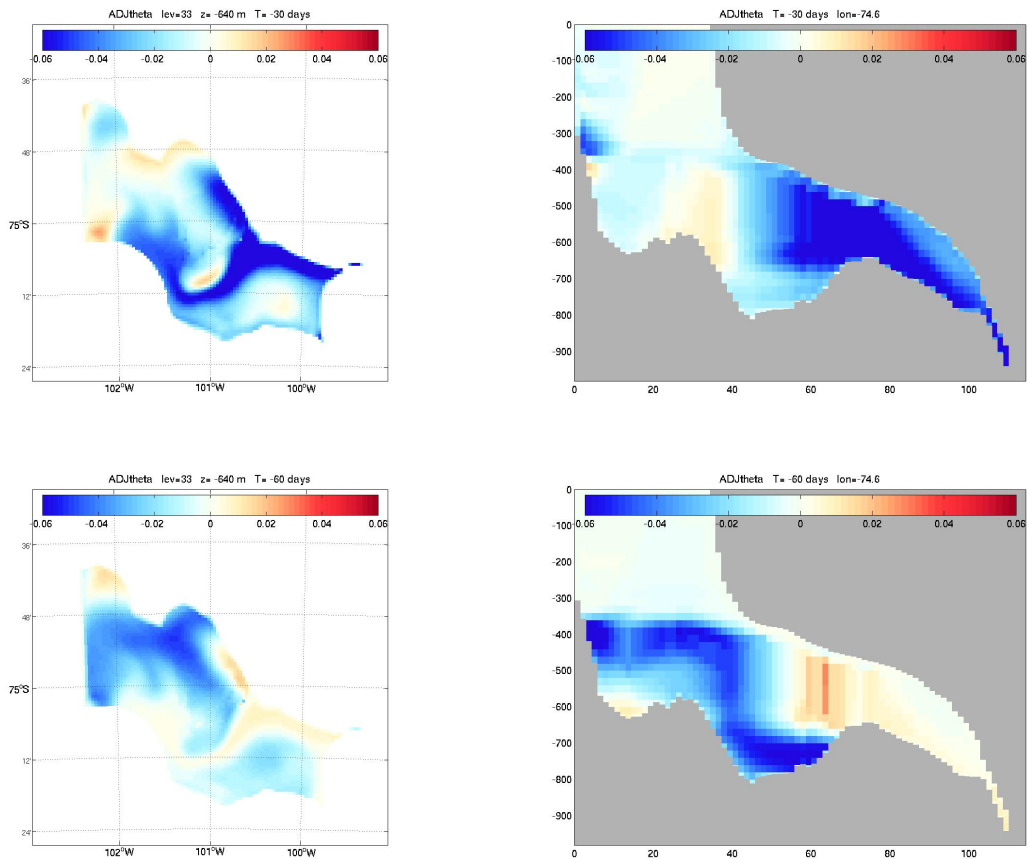


Figure 14: Transient sensitivities, $\delta^*T = (\partial J/\partial T)^T$, of integrated melt rates J under PIIS (from Fig. 11b) to changes in temperature T at times $t = \tau_f - 30$ days (upper row) and -60 days (lower row) prior to computing J . Left panels are horizontal slices at 640 m depth, right panels are vertical slices taken along the dashed line depicted in Fig. 13. Units are in $\text{m}^3 \text{s}^{-1} \text{K}^{-1}$, where $0.1 \text{ m}^3 \text{s}^{-1} \text{K}^{-1} \approx 3 \text{ Mt a}^{-1} \text{K}^{-1} \approx 3 \text{ mm a}^{-1} \text{K}^{-1}$.

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5.2 Longer Duration Estimates

Although the original ECCO estimates were confined to the period beginning in the early 1990s with the improved observational coverage that became available in association with WOCE, the intense interest in decadal scale climate change has led to some estimates of the ocean state emulating the meteorological reanalyses, extending 50 years and longer into the past. Some of these estimates are based essentially on the reanalysis methods already described (e.g., Rosati et al., 1995; Hurlburt et al., 2009), and having all of their known limitations.

Köhl and Stammer (2008), Wang et al. (2010) have pioneered the application of the ECCO least-squares methods to an oceanic state estimate extending back to 1960. Their estimates have

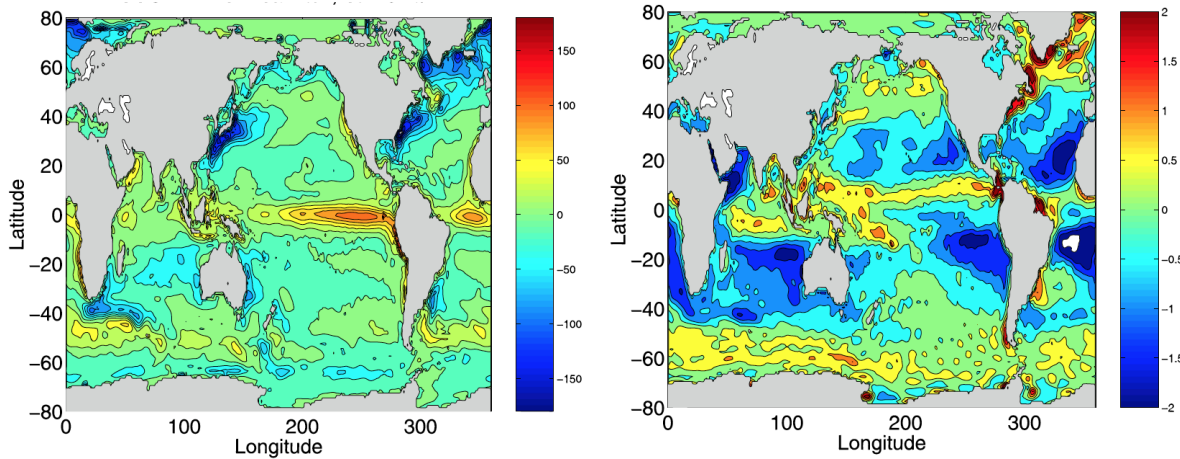


Figure 15: From Stammer et al. (2004) showing an estimate of the multi-year average heat (left, in W/m^2) and fresh water (right, in m/y) transfers between ocean and atmosphere.

{stammer_josey

614 the same virtue as the wider ECCO family of solutions, in satisfying known model equations
 615 of motion and dynamics and with known misfits to all data types. The major problem is the
 616 extreme paucity of data in the ocean preceding the WOCE era; see e.g., Figs. 1 and 2 of Forget
 617 and Wunsch (2007), and the accompanying very limited meteorological forcing observations in
 618 the early days. Note that polar orbiting meteorological satellites do not exist prior to 1979—see
 619 Fig. 2 and Bromwich and Fogt (2004). Useful altimetry appears only at the end of 1992. “Whole
 620 domain” methods such as smoothers or Lagrange multipliers do carry information backwards
 621 in time, and in the estimates for the underconstrained decades prior to about 1992, the gross
 622 properties of the ocean circulation are better determined because of the later, denser, data
 623 sets. But the memory of the upper ocean, which is most prominent e.g., in climate forecasting
 624 attempts, appears to be restricted to a few years, and one expects considerable near-surface
 625 uncertainty to occur even as recently as the 1980s.

626 A preliminary step of assessing the impact of observational assets in constraining the ECCO
 627 solutions has been taken through observing system experiments in the context of short-duration
 628 optimizations during the Argo array period (Heimbach et al., 2009; Zhang et al., 2010). Results
 629 suggest that the impact of altimetry and Argo floats in constraining, e.g. the MOC is drastic,
 630 compared to the pre-WOCE period when only hydrographic sections were available.

631 The published solutions for the interval prior to about 1992 are best regarded as physically
 632 possible, but whose uncertainty estimates, were they known, would surely be very much greater
 633 than they are in the later times, but diminishing as the WOCE-era is approached. These long-
 634 duration estimates, decades into the past, thus present a paradox: if they are quantitatively
 635 useful—other than as examples of *possible* solutions—then the relatively large investment in

636 observation systems the community has made since the early 1990s was unnecessary. If that
637 investment has been necessary, then one cannot readily quantitatively interpret the early es-
638 timates. We leave the subject here as one awaiting the necessary time-dependent uncertainty
639 estimates.

640 **5.3 Short-Duration Estimates**

641 Finding a least-squares fit over 19+ years is computationally very demanding and for some
642 purposes, estimates over shorter time intervals can be useful. In particular, Forget (2010) used
643 the same model and methodology as that of the ECCO Climate State 1° system (Wunsch and
644 Heimbach, 2007), but limited the calculation of three overlapping 18-months periods in the
645 years 2004-2006. In his estimate, the model-data misfit is considerably reduced compared to
646 that in the 16+year solution. The reasons for that better fit are easy to understand from the
647 underlying least-squares methodology: The number of adjustable parameters (the control vector)
648 has the same number of degrees-of-freedom in the initial condition elements as does the decade+
649 calculations, but with many fewer data to fit, and with little time to evolve away from the opening
650 state. (Meteorological elements change over the same time scales in both calculations.) It is
651 much more demanding of a model and its initial condition controls to produce fits to a 16-year
652 evolution than to an 18-month one. Although both calculations have time-scales short compared
653 to oceanic equilibrium times of hundreds to thousands of years, in an 18 month interval little
654 coupling exists between the meteorological controls and the deep data sets—which are then
655 easily fit by the estimated initial state.

656 Solutions of this type are very useful, particularly for upper ocean and regional oceanographic
657 estimates (see the water mass formation rate application in Maze et al., 2009). An important
658 caveat, however, is that one must resist the temptation to regard them as climatologies. They
659 *do* bring us much closer to the ancient oceanographic goal of obtaining a synoptic “snapshot.”

660 **5.4 Global High Resolution Solutions**

661 Ocean modelers have been pursuing ever-higher resolution from the very beginning of ocean
662 modelling and the effort continues. In classical computational fluid dynamics, one sought “nu-
663 merical convergence”: the demonstration that further improvements in resolution did not qual-
664 itatively change the solutions, and preferably that they reproduced known analytical values.
665 Such demonstrations with GCMs are almost non-existent, and thus a very large literature has
666 emerged attempting to demonstrate the utility of “parameterizations”—constructs intended to
667 mimic the behavior of motions smaller than the resolution capability of any particular model.
668 A recent review is by Ringler and Gent (2011). Absent fully-resolved solutions with which to

669 compare the newer parameterizations, the question of their quantitative utility remains open.
670 They do represent clear improvements on older schemes.

671 Despite the parameterization efforts, considerable evidence exists (e.g., Hecht and Smith,
672 2008; Lévy et al., 2010) that qualitative changes take place in GCM solutions when the first
673 baroclinic deformation radius, at least, is fully resolved. From the state estimation point of view,
674 one seeks as much skill as possible in the model—which is meant to represent the fullest possible
675 statement of physical understanding. On the other hand, state estimation, as a curve-fitting
676 procedure, is relatively immune to many of the problems of prediction. In particular, because
677 of the dominant geostrophic balance, its mass transport properties are insensitive to unresolved
678 spatial scales—bottom topographic interference being an exception. In data dense regions, away
679 from boundary currents, one anticipates robust results even at modest resolution.

680 Ultimately, however, the boundary current regions particularly must be resolved (no pa-
681 rameterizations exist for unresolved boundary currents) so as to accurately compute transport
682 properties for quantities such as heat or carbon that depend upon the rendering of the second
683 moments, $\langle C\mathbf{v} \rangle$, where C is any scalar property, and \mathbf{v} is the velocity. Thus a major effort has
684 been devoted to producing global or near-global state estimates from higher resolution models
685 (Menemenlis et al., 2005a, b). The same methodologies used at coarser resolution are also ap-
686 propriate at high resolution—as has been demonstrated in the regional estimates taken up next,
687 but the computational load rapidly escalates with the state and control vector dimensions. Thus
688 available globally constrained models have used reduced data sets, and have been calculated only
689 over comparatively short time intervals (see Table 3).

690 Because of the short-duration, much of the interest in these high resolution models lies with
691 the behavior of the eddy field rather than in the large-scale circulation (e.g., Wortham, 2012).
692 As with ordinary forward modelling, how best to adjust the eddy flux parameterizations when
693 parts of the eddy field have been resolved, is a major unknown.

694 **5.5 Regional Solutions**

695 Because the computational load of high resolution global models is so great, efforts have been
696 made to produce regional estimates, typically embedded in a coarser resolution global system.
697 Embedding, with appropriate open boundary conditions is essential, because so much of the
698 ocean state in any finite region is directly dependent upon the boundary values. Implementing
699 open boundary conditions is technically challenging, particularly where the velocity field is
700 directly involved—with slight barotropic imbalances producing large volume imbalances (Ayoub,
701 2006).

702 Gebbie et al. (2006) discussed estimates in a small region of the North Atlantic, and their

703 results were used to calculate (Gebbie, 2007) the eddy contribution to near-surface subduction
704 processes. In a much-larger region, the Mazloff et al. (2010) Southern Ocean State Estimate
705 (SOSE), was computed initially over the restricted time interval 2005-2006 (now being extended)
706 at $1/6^\circ$ horizontal and 42 vertical-level resolution.

707 **6 The Uncertainty Problem**

708 From the earliest days of least-squares as used by Gauss and Lagrange, it was recognized that
709 an important advantage of the methods is their ability to produce uncertainty estimates for the
710 solutions, generally as covariances about the expected solution or the underlying true solution.
711 The art of calculating those errors in historically large systems (especially in geodesy and orbit
712 estimation—the fields where the method originated) is highly refined. Unhappily, large as those
713 systems are, their dimensions pale in comparison to the state and control vector sizes encoun-
714 tered in the oceanographic problem. This dimensionality issue renders impractical any of the
715 conventional means that are useful at small and medium size. Numerous methods have been
716 proposed, including direct calculation of the coefficients of the normal equations (the matrix \mathbf{A} ,
717 defining any system of simultaneous equations) and inversion or pseudo inversion, of $\mathbf{A}^T \mathbf{A}$ (the
718 Hessian); the indirect calculation of the lowest eigenvalues and eigenvectors of the inverse Hes-
719 sian from algorithmic differentiation (AD) tools; to solutions for the probability density through
720 the Fokker-Planck equation; to the generation of ensembles of solutions. Mostly they have been
721 applied to “toy” problems—somewhat similar to designing a bridge to span the Strait of Gibralt-
722 ar, and then pointing at a local highway bridge as a demonstration of its practicality. Serious
723 efforts, more generally, to calculate the uncertainties of any large model solution are continuing,
724 but when a useful outcome will emerge is unknown at this time.

725 In the interim, we generally have only so-called standard errors, representing the temporal
726 variances about the mean of the estimate (Figs. 4, 5). These are useful and helpful. Sensitivi-
727 ties, derived from the adjoint solutions (e.g., Heimbach et al., 2011; and see Figs. 11, 16), are
728 computationally feasible and need to be more widely used. In the meantime, the quest of ocean
729 and climate modelers and for the state estimation community more specifically, for useful un-
730 derstanding of reliability, remains a central, essential, goal. One should note that conventional
731 ocean GCMs or coupled climate models, run without state estimation are almost never accom-
732 panied by uncertainty estimates—a serious lack—particularly in an era in which “prediction
733 skill” is being claimed.

734 Some authors compare their solutions to those inferred from more conventional methods
735 e.g., transport calculations from box inversions of hydrographic sections. These comparisons are

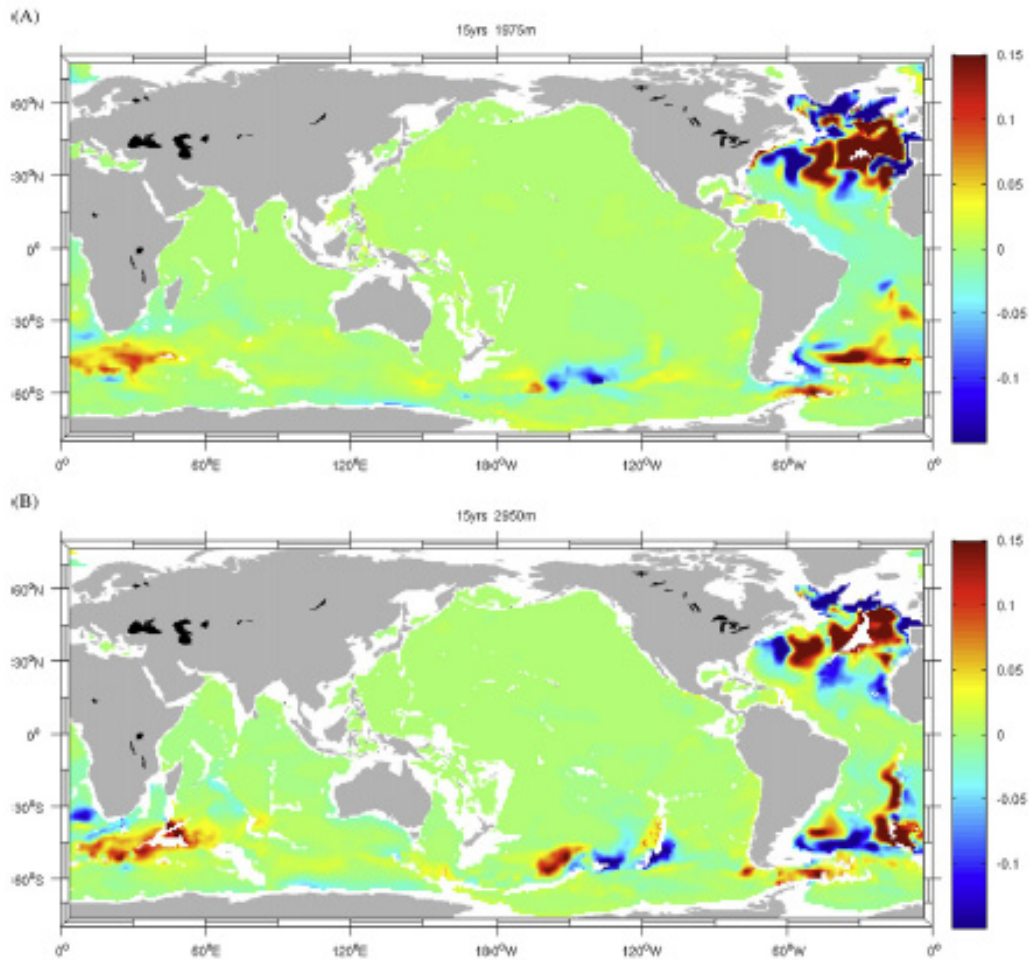


Figure 16: Sensitivities (from Heimbach et al., 2011) of the meridional heat transport across 26°N in the North Atlantic from temperature perturbations at two depths, *15 years earlier*. Top panel is for 1875m, and lower panel is for 2960m.

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736 worthwhile, but it is a major error to treat the hydrographic solutions as if they were true time-
 737 averages or climatologies. It is now possible to compare a state estimate from data obtained
 738 over a short interval (e.g., March 2003) with a state estimate for that time, sampled in the same
 739 way. Differences will appear in the objective function, J . Inevitable discrepancies raise all of the
 740 fundamental questions of allocating errors amongst the data, the model, and external controls.
 741 In the decadal+ prediction problem (not discussed here), by definition there are no data, and
 742 measures of error and skill are far more difficult to obtain. Divergence of IPCC (2007) models
 743 over time (e.g., Schmittner et al., 2005; Stroeve et al., 2007), even where fitted to the historical
 744 observations, is a strong indicator of the fundamental difficulties involved in extrapolating even
 745 systems that appear to give an apparently good fit to historical data, and they are reminiscent

746 of the parable above of fitting cubics to data.

747 **7 Discussion**

748 The history of fluid dynamics generally, and of complex model use in many fields, all support
749 the inference that models unconstrained by data can and do often go wildly wrong (in the wider
750 sense, see e.g., May, 2004; Post and Votta 2005). Readers will recognize the strong point of view
751 taken by the present authors: that models unaccompanied by detailed, direct, comparisons with
752 and constraints by data are best regarded as a kind of science novel.

753 As we go forward collectively, the need to develop methods describing GCM and state es-
754 timate uncertainties is compelling: how else can one combine the quantitative understanding
755 of oceanographic, meteorological and cryospheric physics with the diverse sets of system ob-
756 servations? Such syntheses are the overarching goal of any truly scientific field. Existing state
757 estimates have many known limitations, some of which will be overcome by waiting for the
758 outcome of Moore’s Law over the coming years. Other problems, including the perennial and
759 difficult problem of oceanic mixing and dissipation (Munk and Wunsch, 1998; Wunsch and Fer-
760 rari, 2004) are unlikely simply to vanish with any foreseeable improvement in computer power.
761 Further insight is required.

762 Lack of long-duration, large-scale, observations generates a fundamental knowledge gap.
763 Without the establishment and maintenance of a comprehensive global ocean observing system,
764 which satisfies the stringent requirements for climate research and monitoring, progress over the
765 coming decades will remain limited (Baker et al., 2007).

766 Oceanographers now also directly confront the limits of knowledge of atmospheric processes.
767 Until about 20 years ago, meteorological understanding so greatly exceeded that of the ocean
768 circulation that estimated state errors for the atmosphere were of little concern. The situa-
769 tion has changed emphatically with the global observations starting in WOCE, along with the
770 development of oceanic state estimates.⁸ These estimation systems are better suited for the
771 purposes of climate research than those developed for numerical weather prediction. What is
772 needed now for climate change purposes, are useful state estimation systems including simul-
773 taneously, coupled oceanic, atmospheric, and sea ice physics, and the entirety of the relevant
774 observations in those fields. Thus atmospheric precipitation and evaporation pattern changes
775 can be constrained tightly by changes in the oceanic state. ECCO and related programs have
776 demonstrated how to carry out such recipes. Conventional weather forecast methods are not

⁸The authors have been asked repeatedly at meetings “Why don’t oceanographers adopt the sophisticated methods used by meteorologists?” The shoe, however, is now firmly on the other foot.

777 appropriate, and implementation of a fully coupled state estimation system that will be ongoing
778 is a challenge to governments, university, and research organizations alike. (cf. Bengtsson et al.,
779 2007, who propose a limited step in this direction. Sugiura et al., 2008, and Mochizuki et al.,
780 2009, have made some tentative starts on it.) Surely we must have the capability.

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reanalysis product	net fresh water imbalance [mm/year]		net heat flux imbalance [W/m ²]	
	ocean-only	global	ocean-only	global
NCEP/NCAR-I (1992-2010)	159	62	-0.7	-2.2
NCEP/DOE-II (1992-2004)	740	-	-10	-
ERA-Interim (1992-2010)	199	53	-8.5	-6.4
JRA-25 (1992-2009)	202	70	15.3	10.1
CORE-II (1992-2007)	143	58		

Table 1: Global net heat and freshwater flux imbalances of atmospheric reanalysis products.

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observation	instrument	product/source	area	period	dT
Mean dynamic topography (MDT)	<ul style="list-style-type: none"> • GRACE SM004-GRACE3 • EGM2008/DNSC07 	CLS/GFZ (A.M. Rio) N. Pavlis/Andersen & Knudsen	global global	time-mean	mean
Sea level anomaly (SLA)	<ul style="list-style-type: none"> • TOPEX/POSEIDON • Jason • ERS, ENVISAT • GFO 	NOAA/RADS & PO.DAAC NOAA/RADS & PO.DAAC NOAA/RADS & PO.DAAC NOAA/RADS & PO.DAAC	65°N/S 82°N/S 65°N/S 65°N/S	1993 – 2005 2001 – 2011 1992 – 2011 2001 - 2008	daily daily daily daily
SST	<ul style="list-style-type: none"> • blended, AVHRR (O/I) • TRMM/TMI • AMSR-E (MODIS/Aqua) 	Reynolds & Smith GHRSSST GHRSSST	Global 40°N/S Global	1992 - 2011 1998 - 2004 2001 - 2011	monthly daily daily
SSS	Various in-situ	WOA09 surface	Global	climatology	monthly
In-situ T, S	<ul style="list-style-type: none"> • Argo, P-Alace • XBT • CTD • SEaOS • TOGA/TAO, Pirata 	lfremer D. Behringer (NCEP) various SMRU & BAS (UK) PMEL/NOAA	“global” “gobal” sections SO Tropics	1992 – 2011 1992 – 2011 1992 – 2011 2004 – 2010 1992 – 2011	daily daily daily daily daily
Mooring velocities	<ul style="list-style-type: none"> • TOGA/TAO, Pirata • Florida Straits 	PMEL/NOAA NOAA/AOML	Trop. Pac. N. Atl.	1992 – 2006 1992 – 2011	daily daily
Climatological T,S	<ul style="list-style-type: none"> • WOA09 • OCCA 	WOA09 Forget, 2010	“global” “global”	1950 - 2000 1950 - 2002	mean mean
sea ice cover	<ul style="list-style-type: none"> • satellite passive microwave radiometry 	NSIDC (bootstrap)	Arcitic, SO	1992 - 2011	daily
Wind stress	QuickScat	<ul style="list-style-type: none"> • NASA (Bourassa) • SCOW (Risien & Chelton) 	global	1999 – 2009 climatolggy	daily monthly
Tide gauge SSH	Tide gauges	NBDC/NOAA	sparse	1992 - 2006	monthly
Flux constraints	from ERA-Interim, JRA-25, NCEP, CORE-2 variances	Various	global	1992 - 2011	2-day to 14-day
Balance constraints			global	1992 - 2011	mean
bathymetry		Smith & Sandwell, ETOPO5	global	-	-

Table 2: Data used in the ECCO global 1° resolution state estimates until about 2011.

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label & version	hor./ver. grid	domain	duration	scope	reference
ECCO-Production Sustained production of decadal climate state estimates (former ECCO-GODAE)					
ver.0 (ECCO-MIT)	2° / 22	80° N/S	1992—1997	first ECCO product – proof of feasibility	Stammer et al. (2002/04)
ver.1 (ECCO-SIO)	1° / 23	80° N/S	1992—2002	begin of 1° sustained production	Köhl et al. (2007)
ver.2 (ECCO-GODAE)	1° / 23	80° N/S	1992—2004	air-sea flux constraints for sea level studies	Wunsch & Heimbach (2006/07)
(OCCA)	1° / 50	80° N/S	2004/5/6/7	Atlas from one-year “synoptic snapshots”	Forget (2010)
(GECCO)	1° / 23	80° N/S	1951—2000	50-year solution covering NCEP/NCAR period	Köhl and Stammer (2008a/b)
ver.3.0 (ECCO-GODAE)	1° / 23	80° N/S	1992—2007	switch to atmos. state controls and sea ice	Wunsch & Heimbach (2009)
revision 1 (ver.3.1)	1° / 23	80° N/S	1992—2010	updates to ver.3.0	Fukumori et al. (in prep.)
ver.4	1° / 50	global	1992—2010	first full-global estimate incl. Arctic	Forget & Heimbach (in prep.)
ECCO-ICES Ocean-Ice Interactions in Earth System Models (former ECCO2)					
ver.1 (CS510 GF)	18 km / 50	global	1992—2002/10	Green’s function optim., of eddying model	Menemenlis et al. (2005)
ver.2 (CS510 adjoint)	18 km / 50	global	2004-05/2009-10	adjoint-based global eddying 1-year optim.	Menemenlis et al. (in prep.)
ECCO-JPL near real-time filter & reduced-space smoother					
ver.1 (KF)	1° / 46	80° N/S	1992—present	near-real time Kalman Filter (KF) assimilation	Fukumori et al. (1999)
ver.2 (RTS)	1° / 46	80° N/S	1992—present	smoother update of KF solution	Fukumori (2002)
Regional Efforts (/*/ denote ongoing efforts)					
Southern Ocean (SOSE) /*/	1/6° / 42	25°-80°S	2005—2009	eddy-permitting SO State Estimate	Mazloff et al. (2010)
ECC2 Arctic & ASTE /*/	18 & 4 km / 50	Arctic & SPG	1992—2009	Arctic/subpolar gyre ocean-sea ice estimate	Nguyen et al. (2011/12)
North Atlantic	1° / 23	25°-80°N	1993	experimental 2° vs. 1° nesting	Ayoub (2006)
Subtropical Atlantic	1/6° / 42	...	1992/93	experimental 1° vs. 1/6° nesting	Gebbie et al. (2006)
Tropical Pacific		experimental 1° vs. 1/3° nesting	Hoteit et al. (2006/2010)
Labrador Sea & Baffin Bay	1996/97	first full coupled ocean-sea ice estimate	Fenty et al. (2012a/b)

Table 3: Published ECCO family state estimates, divided roughly into categories. The global decade+ estimates are labelled as “ECCO-Production”, while others are either regional, or experimental.

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